

SOLVING FOR THE LOW-RANK TENSOR COMPONENTS OF THE WAVE FUNCTION IN SCATTERING PROBLEMS WITH MULTIPLE IONIZATION



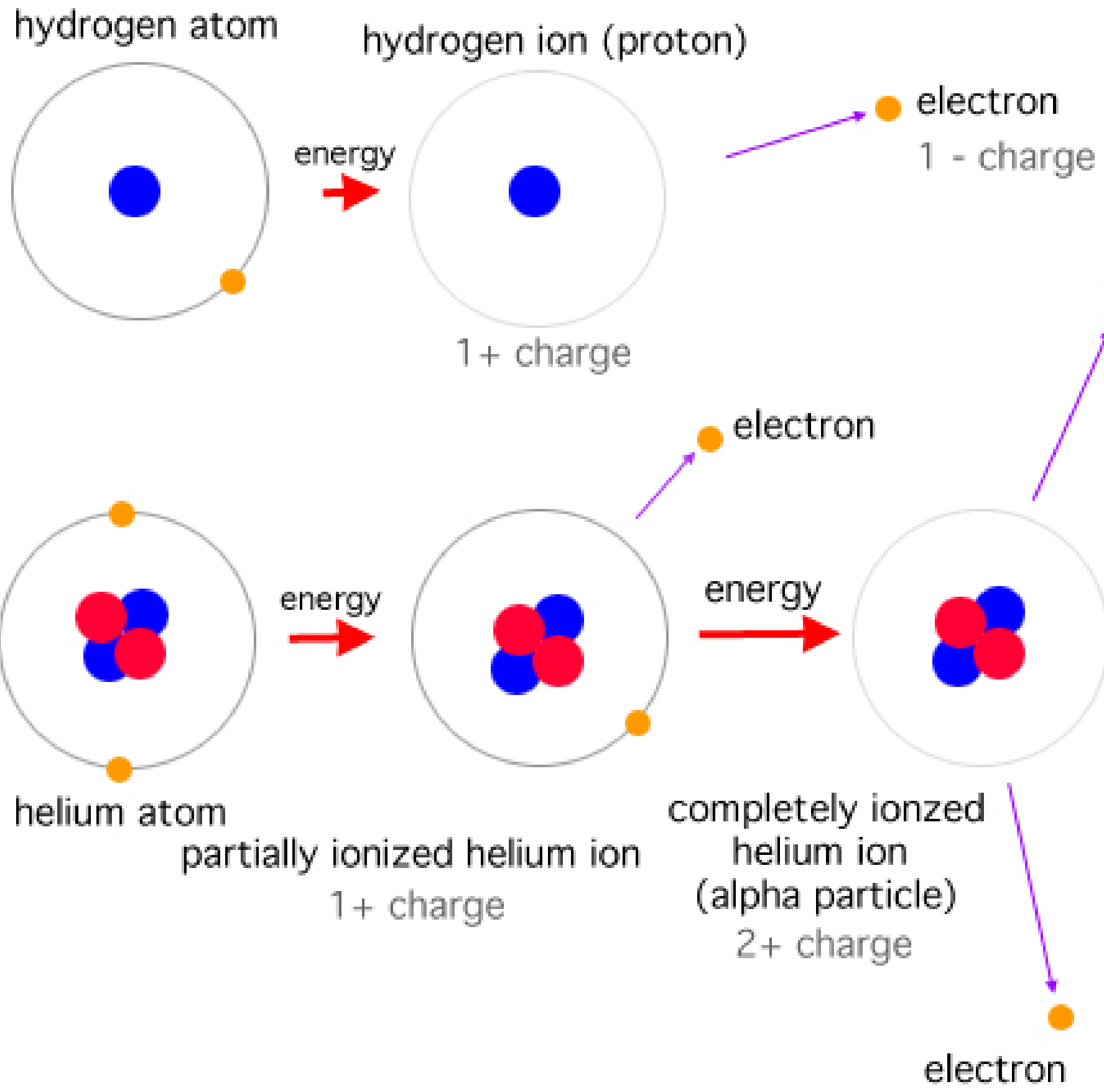
Jacob Snoeijer (jacob.snoeijer@uantwerpen.be) and Wim Vanroose



University of Antwerp, Department of Mathematics

Ionization & scattering problems

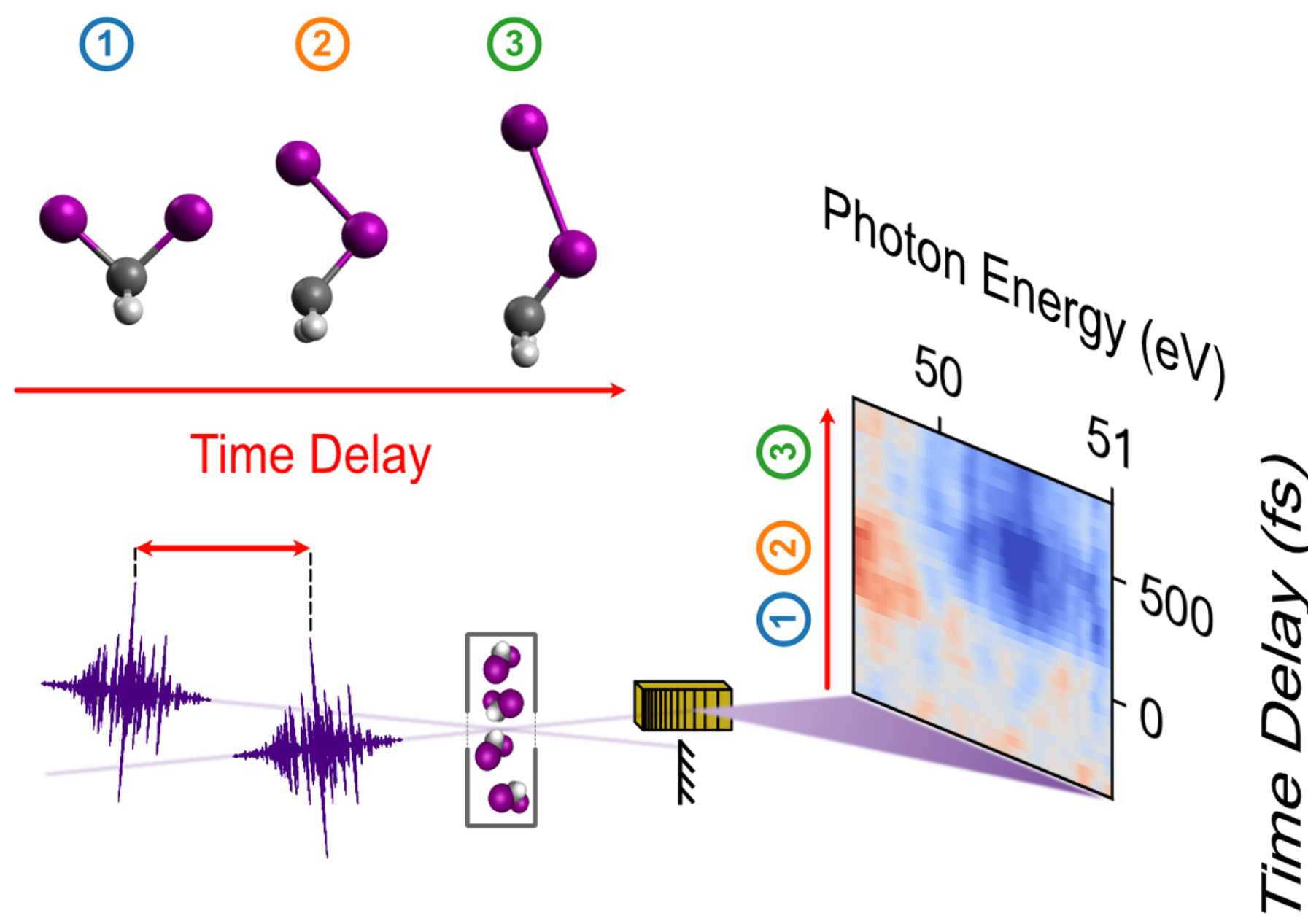
Information about microscopic systems comes from experiments where electrons scatter from an object.



Molecular movie

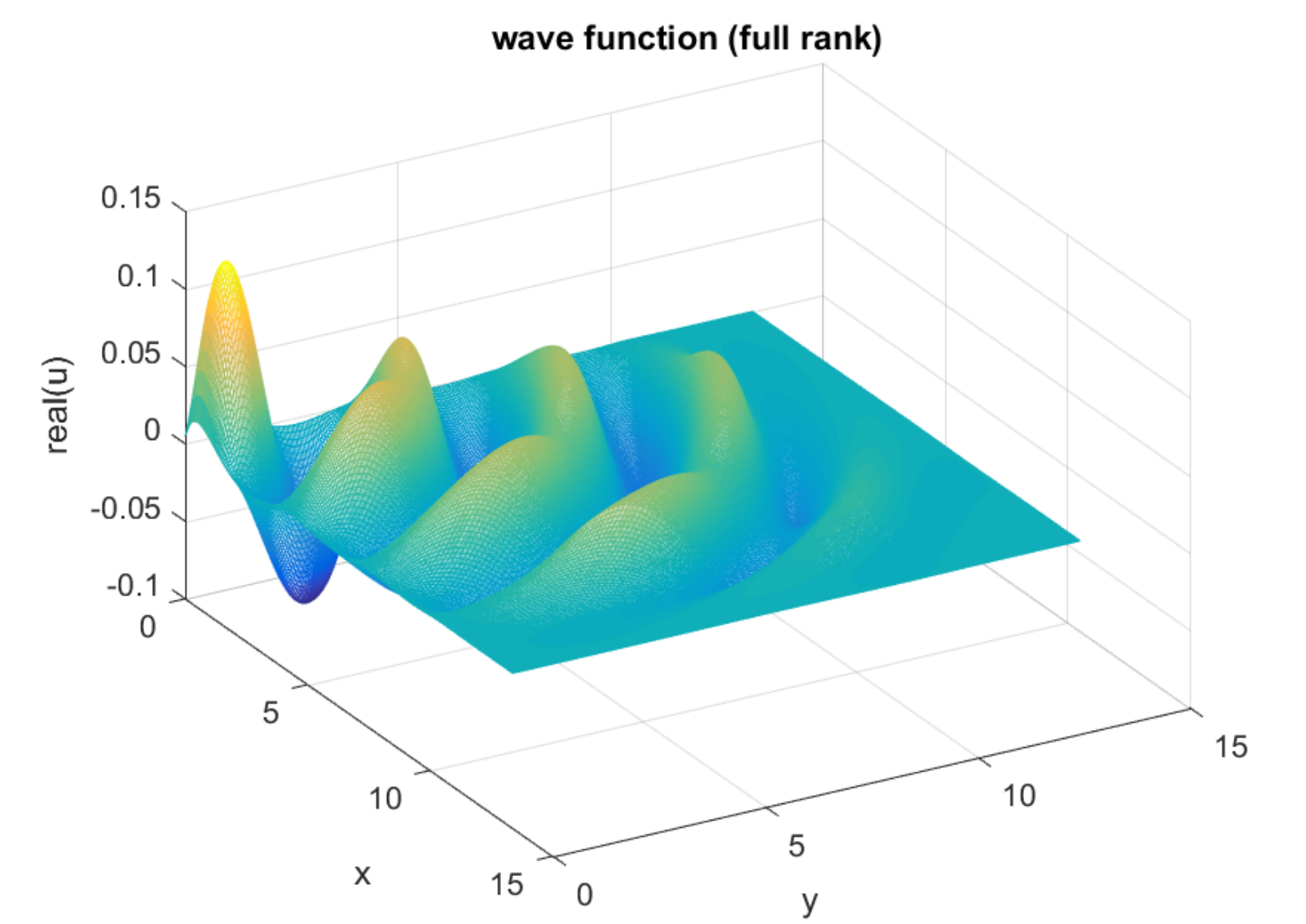
Electrons are emitted along angles that can be measured. This results in a probability distribution, called the **far field map**. This corresponds to the amplitude of the wave in the emitting direction.

In new experiments, based on a far field map, the original state of molecules will be reconstructed.



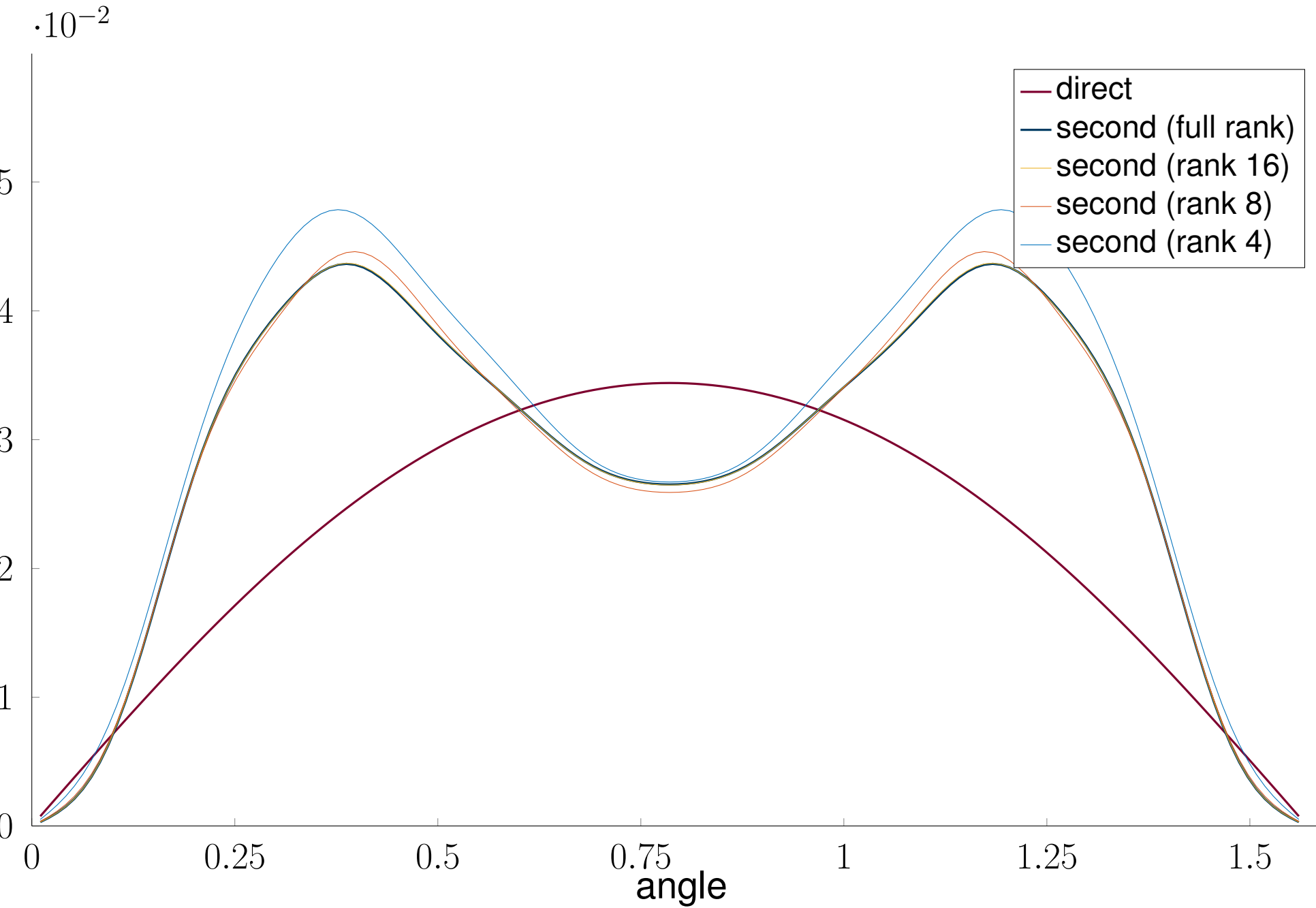
Wave function

- 2D Helmholtz equation with $(x, y) \in (0, 10)^2$ and extended with ECS, $E = 4$, $\omega^2(x, y) = \exp(-|(x-y)^4|) - E$, $f(x, y) = \exp(-x^2 - y^2)$.
- Discretization: FD grid of 200 grid points and ECS layer with 67 points rotated over $\pi/6$.

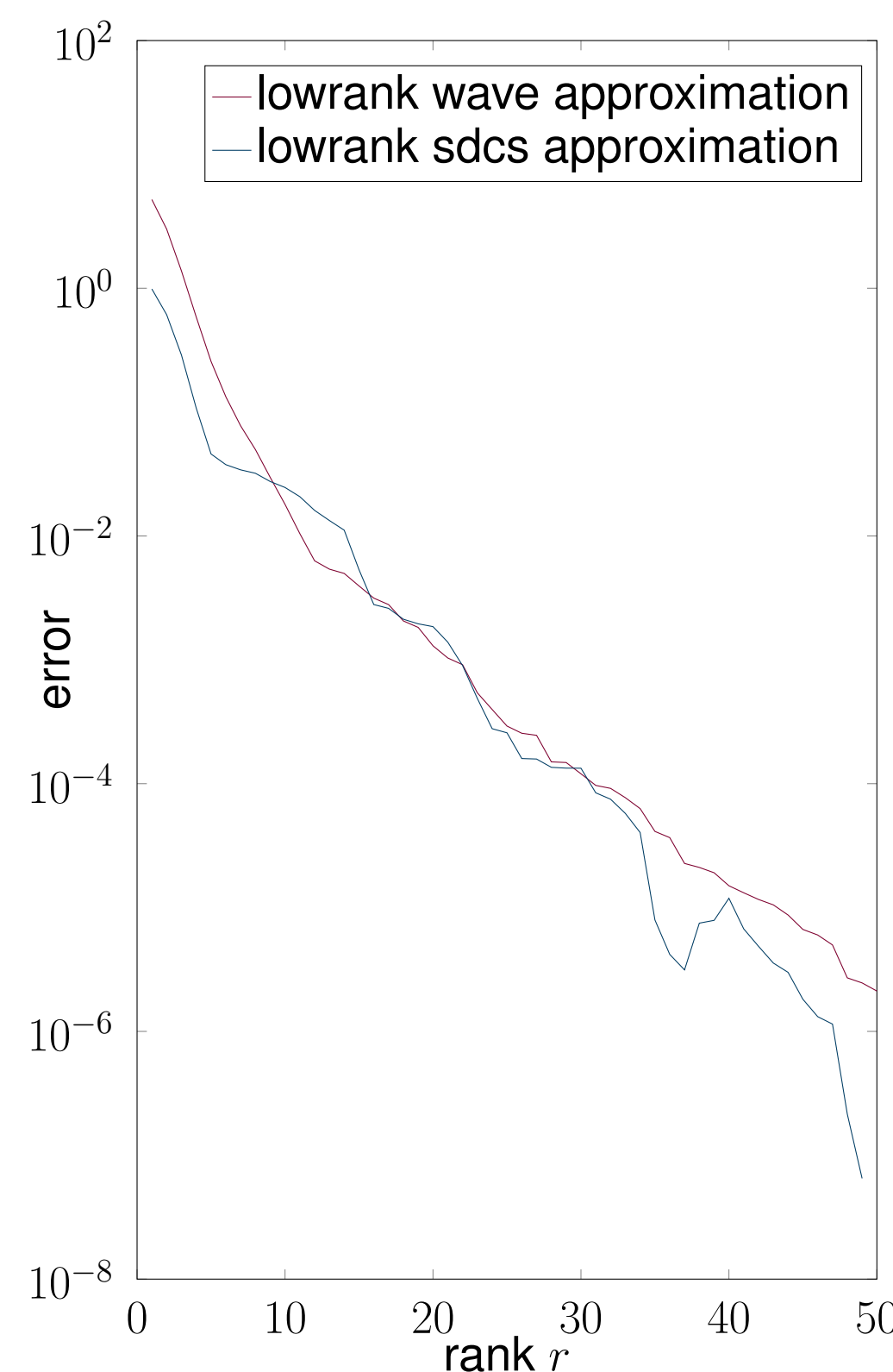


Differential cross section

$$\frac{d\sigma}{d\Omega}(\theta) = \int_0^b \int_0^b \sin(\omega_0 \sin \theta x) \sin(\omega_0 \cos \theta y) [f + \omega_0^2 \chi u_{sc}] dx dy$$



Error SDCS



Observation (2D)

To get a good approximation for the SDCS a **low rank** approximation to the Helmholtz equation is sufficient.

Discretization of Helmholtz equation on a 2D cartesian grid with $n_x \times n_y$ unknowns leads to the following matrix equation:

$$D_{xx}A + AD_{yy}^T - W \circ A = F \quad (1)$$

where A , F and W are $n_x \times n_y$ matrices that describe the unknown solution $u_{sc}(x, y)$, f and ω^2 on the discretized mesh.

Assuming that A has low rank this matrix can be written as $A = UV^H$ with $U \in \mathbb{C}^{n_x \times r}$ and $V \in \mathbb{C}^{n_y \times r}$ both with orthogonal columns and $r \ll \min(n_x, n_y)$ is the rank of matrix A :

$$D_{xx}UV^H + UV^H D_{yy}^T - \omega^2 UV^H = F \quad (2)$$

Derive equations to solve for the factors U and V :

$$\begin{aligned} D_{xx}U + UV^H D_{yy}^T V - \omega^2 U &= FV \\ U^H D_{xx} UV^H + V^H D_{yy}^T V - \omega^2 V^H &= U^H F \end{aligned} \quad (3)$$

The algorithm (2D)

We have update equations for U and V :

$$\begin{aligned} [(I \otimes (D_{xx} - \omega^2 I)) + (V^T D_{yy} \bar{V} \otimes I)] \text{vec}[U] &= \text{vec}[FV] \\ [(I \otimes U^H D_{xx} U) + ((D_{yy} - \omega^2 I) \otimes I)] \text{vec}[V^H] &= \text{vec}[U^H F] \end{aligned}$$

Given: initial guess $V_0 \in \mathbb{C}^{n_y \times r}$.

$[V, R] = \text{qr}[V_0]$;

while not converged do

Solve for U ;

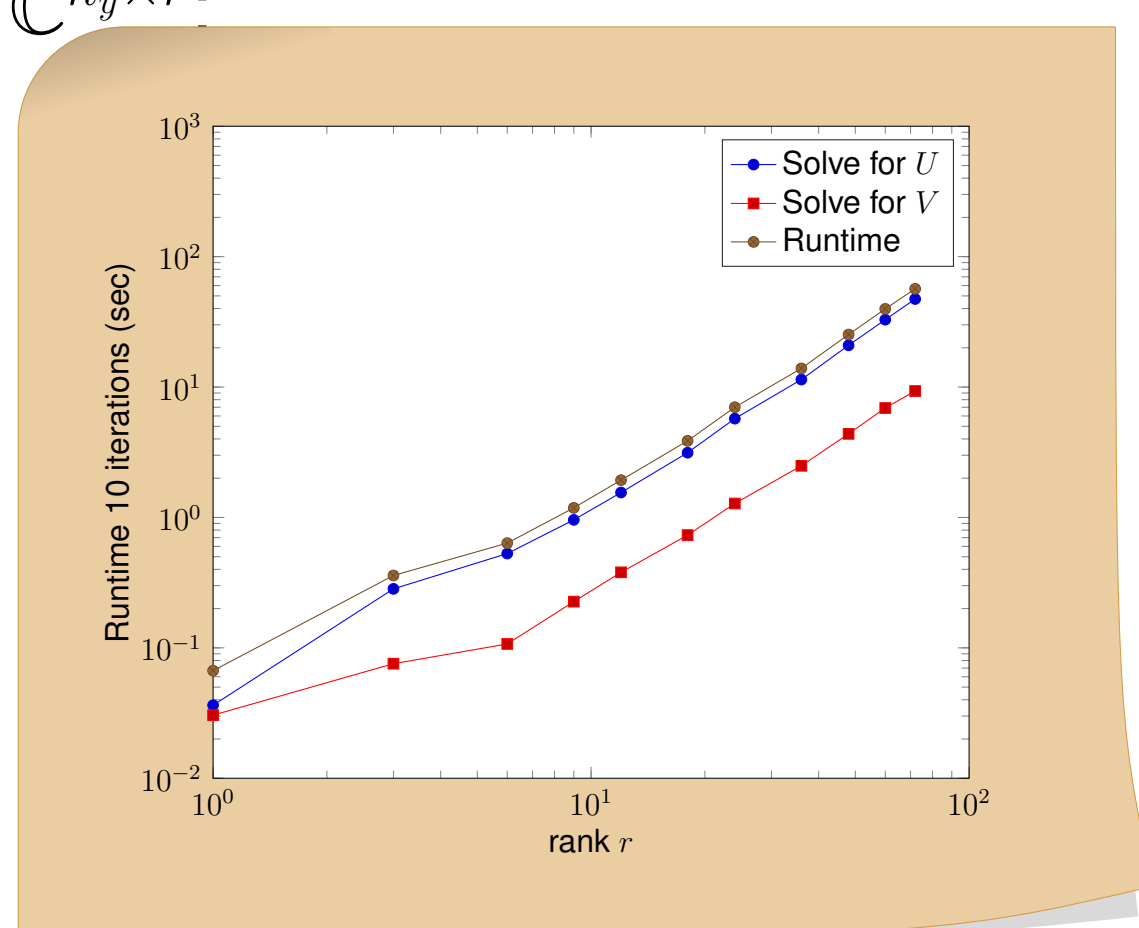
$[U, \tilde{R}] = \text{qr}[U]$;

Solve for V^H ;

$[V, R] = \text{qr}[V]$;

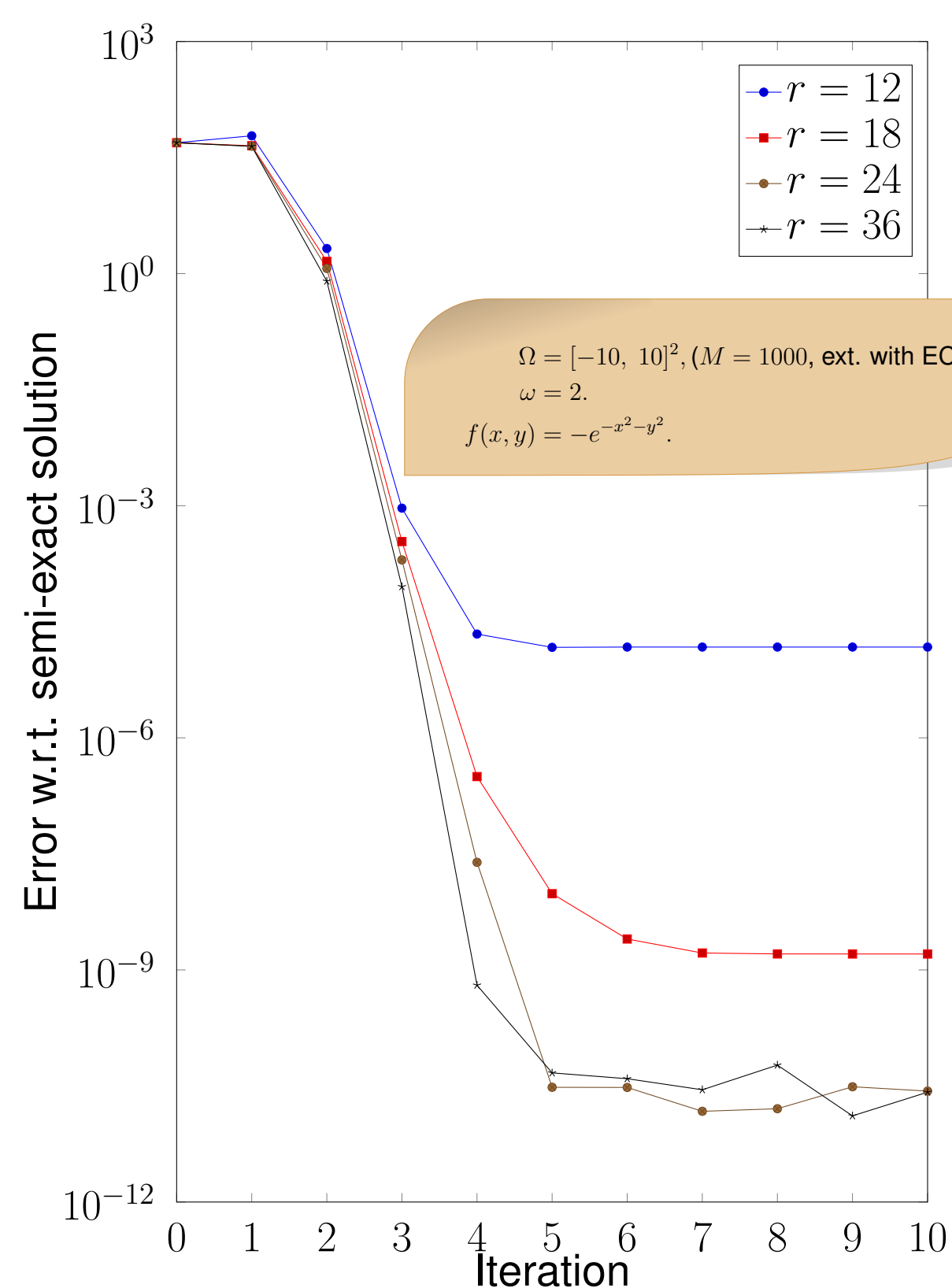
end

$A = UR^H V^H$;



- Solving for U , V can be interpreted as a **projection operator** applied on residual.
- The successive application of these projection operators doesn't result in a new projection operator. But it can be seen as a **perturbation** of a projection operator.
- This perturbation is typical **small and bounded**, which leads to convergence.

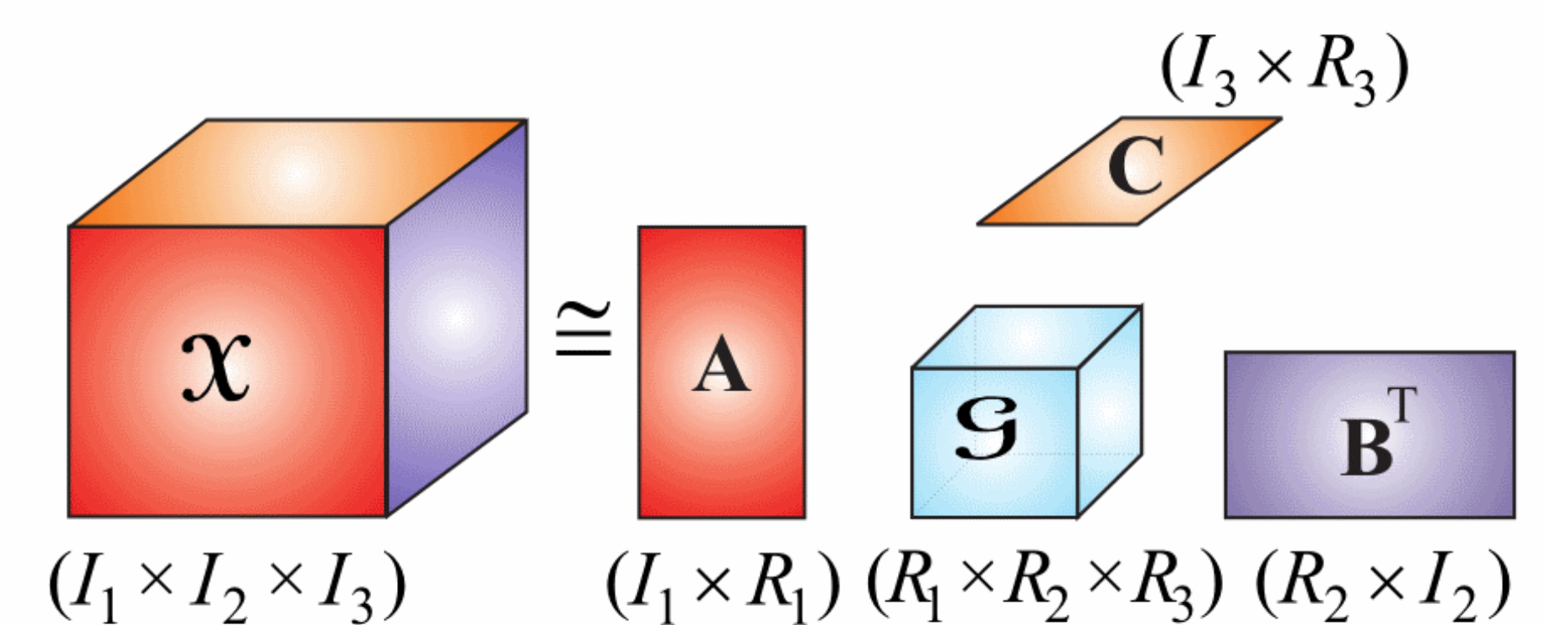
Errors per iteration



Tucker tensor decomposition

A matrix SVD is generalized by a **Tucker tensor decomposition** which represents an d -th order tensor $\mathcal{X} \in \mathbb{R}^{I_1 \times I_2 \times \dots \times I_d}$ as a multilinear transformation of a dense core tensor $\mathcal{G} \in \mathbb{R}^{R_1 \times R_2 \times \dots \times R_d}$ by orthonormal matrices $U_n = U^{(n)} \in \mathbb{R}^{I_n \times R_n}$:

$$\begin{aligned} \mathcal{X} &= \mathcal{G} \times_1 U^{(1)} \times_2 U^{(2)} \dots \times_d U^{(d)} \\ &= \sum_{r_1=1}^{R_1} \sum_{r_2=1}^{R_2} \dots \sum_{r_d=1}^{R_d} g_{r_1 r_2 \dots r_d} (u_{r_1}^{(1)} \circ u_{r_2}^{(2)} \circ \dots \circ u_{r_d}^{(d)}) \end{aligned} \quad (4)$$



Helmholz on tensors

Application of Helmholtz opr. \mathcal{L} on \mathcal{M} :

$$\begin{aligned} \mathcal{L}\mathcal{M} = \mathcal{F} &= \mathcal{G} \times_1 D_{xx} U_1 \times_2 U_2 \times_3 U_3 \\ &+ \mathcal{G} \times_1 U_1 \times_2 D_{yy} U_2 \times_3 U_3 \\ &+ \mathcal{G} \times_1 U_1 \times_2 U_2 \times_3 D_{zz} U_3 \\ &- \mathcal{W} \circ (\mathcal{G} \times_1 U_1 \times_2 U_2 \times_3 U_3) \end{aligned}$$

where $U_i^H U_i = I$, for $i = 1, 2, 3$.

The algorithm (3D)

To derive an update equation for U_1 multiply with U_2 and U_3 in the second and third dimension, respectively:

$$\mathcal{L}\mathcal{M} \times_2 U_2^H \times_3 U_3^H = \mathcal{F} \times_2 U_2^H \times_3 U_3^H$$

Unfolding in first mode leads to a matrix equation for $\bar{U}_1 \mathcal{G}_{(1)}$.

Conclusions

- In certain cases the solution to Helmholtz can be described by a **limited number of singular values**. Low rank approximations are sufficient and **reduce computational cost** to construct a far field map.
- The presented algorithm has **good contractivity properties**. In practice only a **small number of iterations** are needed (also for low rank space dependent wave numbers).
- Using Tucker decompositions the presented algorithm can be **extended to nD**.