

# Hybrid Numerical-Asymptotic Methods for High Frequency Scattering Problems

**Simon Chandler-Wilde**  
University of Reading, UK

[www.reading.ac.uk/~sms03snc](http://www.reading.ac.uk/~sms03snc)

With: **Steve Langdon**, **Andrea Moiola** (Reading), **Ivan Graham**,  
**Euan Spence** (Bath), **Dave Hewett** (Oxford), **Valery Smyshlyayev**,  
**Timo Betcke** (UCL), **Marko Lindner** (TU HH), **Peter Monk** (Delaware)  
PhDs **Andrew Gibbs**, **Sam Groth**, Charlotta Howarth, Ashley Twigger

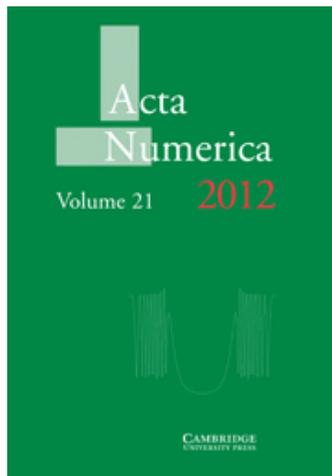
Funding: • EPSRC project(s) across Bath/Reading/UCL with  
**BAE Systems**, **Institute of Cancer Research**, **Met Office**,  
**Schlumberger Cambridge Research** as project partners.  
• 4 NERC/EPSRC CASE Studentships at Bath & Reading  
• EPSRC/Swiss NSF Fellowships for **Andrea**, **Euan**, **Timo**

**Woudschoten, October 2014**

## Context

My talks apply (particularly) to **acoustic waves**.

My talks concern new **numerical-asymptotic methods** for **high frequency** wave scattering based on **boundary integral equations**, that combine **numerical analysis** with **high frequency asymptotics**, see



C-W, Graham, Langdon, Spence *Acta Numerica* 21 (2012), 89–305.

The first talk was largely **HF asymptotics** – in this talk we come to **numerical methods** and their analysis!

## Overview

1. Green's Representation Theorem and **boundary integral equations**  
(for Helmholtz)

2. A Case study for **numerical-asymptotic** methods: **the thin screen**

Step A. Represent the unknown as sum of products of **known oscillatory**  
and **unknown non-oscillatory** functions using **GO/GTD**.

Step B. Decide on the approximation space - combine **HF asymptotics**  
with *hp*-**approximation theory**

Step C. Implement it and see that (we hope) the cost is  $O(1)$  as  $k \rightarrow \infty$ !

Step D. Try to prove this by theorems about the  $k$ -dependence of  
everything!

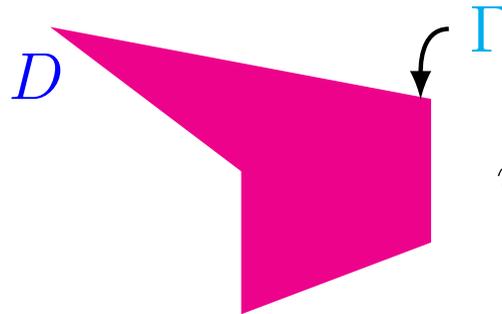
3. **Other geometries and 3D**

**1. GREEN'S REPRESENTATION THEOREM AND INTEGRAL EQUATIONS FOR HELMHOLTZ**

## Green's Representation Theorem: slide from Talk 1

$\mathcal{N} \rightarrow u^{\text{inc}}$

$$\Delta u + k^2 u = 0$$



$u - u^{\text{inc}}$  satisfies S.R.C.

### Theorem

$$u(x) = u^{\text{inc}}(x) + \int_{\Gamma} \left( \frac{\partial u}{\partial n}(y) \Phi(x, y) - u(y) \frac{\partial \Phi(x, y)}{\partial n(y)} \right) ds(y), \quad x \in D,$$

where

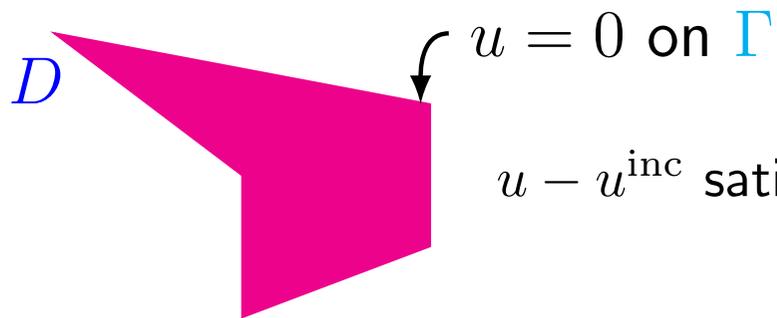
$$\Phi(x, y) = \begin{cases} \frac{i}{4} H_0^{(1)}(k|x-y|) & (2D), \\ \frac{1}{4\pi} \frac{e^{ik|x-y|}}{|x-y|}, & (3D). \end{cases}$$

## Green's Representation Theorem

$\mathcal{N} \rightarrow u^{\text{inc}}$

$$\Delta u + k^2 u = 0$$

$D$



$u = 0$  on  $\Gamma$

$u - u^{\text{inc}}$  satisfies S.R.C.

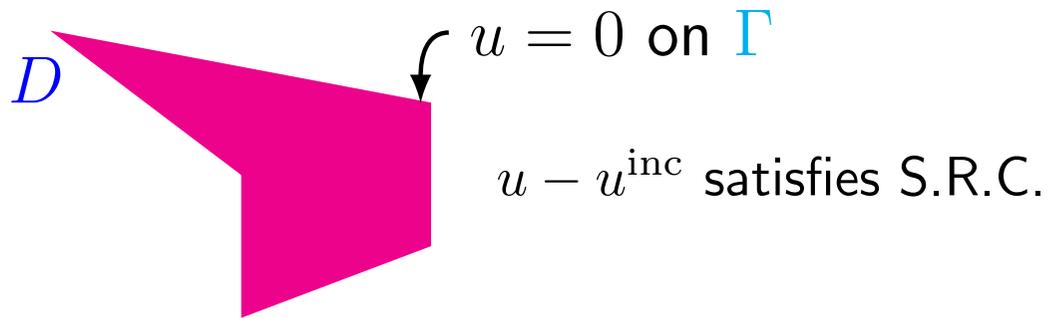
### Theorem

$$u(x) = u^{\text{inc}}(x) + \int_{\Gamma} \left( \frac{\partial u}{\partial n}(y) \Phi(x, y) - u(y) \frac{\partial \Phi(x, y)}{\partial n(y)} \right) ds(y), \quad x \in D.$$

## Green's Representation Theorem

$\mathcal{W} \rightarrow u^{\text{inc}}$

$$\Delta u + k^2 u = 0$$

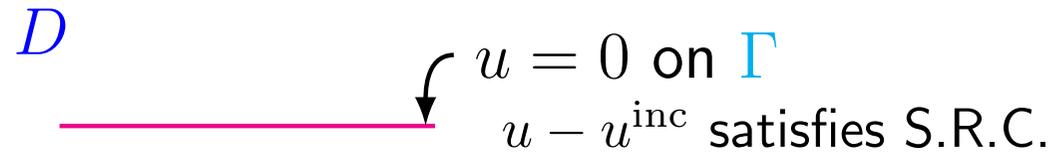


### Theorem

$$u(x) = u^{\text{inc}}(x) + \int_{\Gamma} \frac{\partial u}{\partial n}(y) \Phi(x, y) ds(y), \quad x \in D.$$

## Green's Representation Theorem for a Thin Screen

$$\mathcal{W} \rightarrow u^{\text{inc}} \quad \Delta u + k^2 u = 0$$

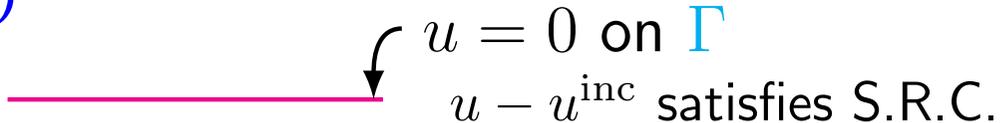


### Theorem

$$u(x) = u^{\text{inc}}(x) + \underbrace{\int_{\Gamma} \left[ \frac{\partial u}{\partial n} \right]}_{\text{jump}}(y) \Phi(x, y) ds(y), \quad x \in D.$$

## Green's Representation Theorem and BIE for a Thin Screen

$$\mathcal{N} \rightarrow u^{\text{inc}} \quad \Delta u + k^2 u = 0$$

$D$   

 $u = 0$  on  $\Gamma$   
 $u - u^{\text{inc}}$  satisfies S.R.C.

### Theorem

$$u(x) = u^{\text{inc}}(x) + \int_{\Gamma} \left[ \frac{\partial u}{\partial n} \right] (y) \Phi(x, y) ds(y), \quad x \in D.$$

Further (letting  $x \rightarrow \Gamma$ ),

$$0 = u^{\text{inc}}(x) + \int_{\Gamma} \left[ \frac{\partial u}{\partial n} \right] (y) \Phi(x, y) ds(y), \quad x \in \Gamma.$$

## Green's Representation Theorem and BIE for a Thin Screen

$$\mathcal{W} \rightarrow u^{\text{inc}} \quad \Delta u + k^2 u = 0$$

**Theorem**

$D$    $u = 0$  on  $\Gamma$   
 $u - u^{\text{inc}}$  satisfies S.R.C.

$$u(x) = u^{\text{inc}}(x) + \int_{\Gamma} \left[ \frac{\partial u}{\partial n} \right] (y) \Phi(x, y) ds(y), \quad x \in D,$$

and

$$0 = u^{\text{inc}}(x) + \int_{\Gamma} \left[ \frac{\partial u}{\partial n} \right] (y) \Phi(x, y) ds(y), \quad x \in \Gamma,$$

in operator notation

$$S_k \left[ \frac{\partial u}{\partial n} \right] = -u^{\text{inc}}|_{\Gamma} \quad \text{(BIE)}$$

## **2. A CASE STUDY FOR NUMERICAL-ASYMPTOTIC METHODS: THE THIN SCREEN**

## A Case Study for Numerical-Asymptotic Methods

$$\mathcal{W} \rightarrow u^{\text{inc}} \quad \Delta u + k^2 u = 0$$

$D$

$u = 0$  on  $\Gamma$   
 $u - u^{\text{inc}}$  satisfies S.R.C.

**Theorem**

$$u(x) = u^{\text{inc}}(x) + \int_{\Gamma} \left[ \frac{\partial u}{\partial n} \right] (y) \Phi(x, y) ds(y), \quad x \in D,$$

and

$$0 = u^{\text{inc}}(x) + \int_{\Gamma} \left[ \frac{\partial u}{\partial n} \right] (y) \Phi(x, y) ds(y), \quad x \in \Gamma,$$

in operator notation

$$S_k \left[ \frac{\partial u}{\partial n} \right] = -u^{\text{inc}}|_{\Gamma} \quad \text{(BIE)}$$

## A Case Study for Numerical-Asymptotic Methods

$$\mathcal{N} \rightsquigarrow u^{\text{inc}} \quad \Delta u + k^2 u = 0$$

$$D \quad \text{---} \quad \begin{array}{l} u = 0 \text{ on } \Gamma \\ u - u^{\text{inc}} \text{ satisfies S.R.C.} \end{array}$$

We will solve

$$0 = u^{\text{inc}}(x) + \int_{\Gamma} \left[ \frac{\partial u}{\partial n} \right] (y) \Phi(x, y) ds(y), \quad x \in \Gamma,$$

in operator notation

$$S_k \left[ \frac{\partial u}{\partial n} \right] = -u^{\text{inc}}|_{\Gamma}, \quad \text{(BIE)}$$

by a **Galerkin BEM** representing  $\left[ \frac{\partial u}{\partial n} \right]$  as a sum of products of **known oscillatory functions** and **unknown non-oscillatory functions**.

## A Case Study for Numerical-Asymptotic Methods

$\mathcal{N} \rightarrow u^{\text{inc}}$

$$\Delta u + k^2 u = 0$$

$D$

$u = 0$  on  $\Gamma$   
 $u - u^{\text{inc}}$  satisfies S.R.C.

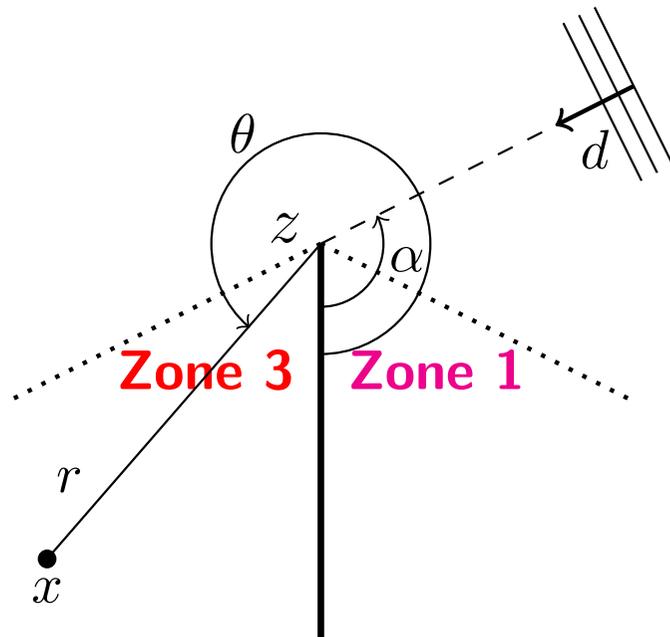
We will solve the BIE by a Galerkin BEM representing  $\left[\frac{\partial u}{\partial n}\right]$  as a sum of products of **known oscillatory functions** and **unknown non-oscillatory functions**.

The steps are:

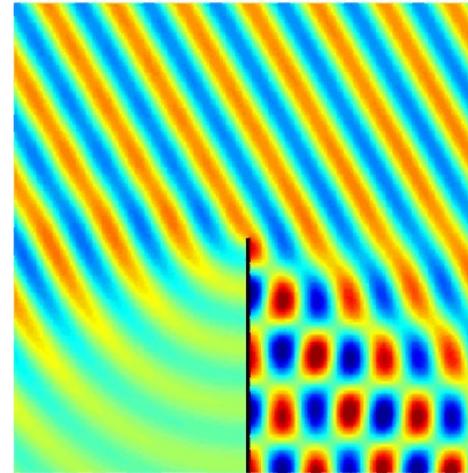
- Work out what this representation for  $\left[\frac{\partial u}{\partial n}\right]$  is - use **GTD**.
- Decide on the approximation space - use **HF asymptotics** and ***hp*-approximation** theory
- Implement it and see that the cost is  $O(1)$  as  $k \rightarrow \infty$ !
- Try to prove this!

**STEP A. REPRESENT THE UNKNOWN AS SUM OF PRODUCTS OF KNOWN OSCILLATORY AND UNKNOWN NON-OSCILLATORY FUNCTIONS USING GO/GTD.**

## The GTD Approach to Diffraction: Recap from Talk 1



$$u^{\text{inc}}(x) = e^{ikx \cdot d}$$



$$\Re u(x)$$

$$u(x) \approx \begin{cases} u^{\text{inc}}(x) + u^{\text{ref}}(x) + u^{\text{d}}(x), & x \text{ in Zone 1} \\ u^{\text{d}}(x), & x \text{ in Zone 3} \end{cases}$$

where  $u^{\text{d}}(x) = u^{\text{inc}}(z) \mathcal{D}(\theta, \alpha) \frac{e^{ikr}}{\sqrt{kr}}$ ,  $\mathcal{D}$  the **diffraction coefficient**.

$\mathcal{N}_{\rightarrow} u^{\text{inc}}$  $D$ 

$u = 0$  on  $\Gamma$

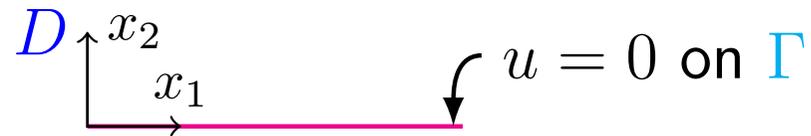
We will solve the BIE by a Galerkin BEM representing  $\left[\frac{\partial u}{\partial n}\right]$  as a sum of products of **known oscillatory functions** and **unknown non-oscillatory functions**.

Step A. Work out what this representation for  $\left[\frac{\partial u}{\partial n}\right]$  is - use **GTD**.

Above the screen (see the GO example yesterday)

$$\frac{\partial u}{\partial n} \approx 2 \frac{\partial u^{\text{inc}}}{\partial n}$$

$\mathcal{N} \rightarrow u^{\text{inc}}$



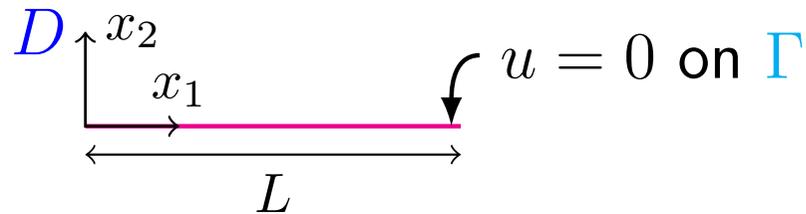
We will solve the BIE by a Galerkin BEM representing  $\left[\frac{\partial u}{\partial n}\right]$  as a sum of products of **known oscillatory functions** and **unknown non-oscillatory functions**.

Step A. Work out what this representation for  $\left[\frac{\partial u}{\partial n}\right]$  is - use **GTD**.

Above the screen ... and adding in the GTD terms ...

$$\frac{\partial u}{\partial n}(x) \approx 2 \frac{\partial u^{\text{inc}}}{\partial n}(x) + c_+ \frac{e^{ikx_1}}{\sqrt{kx_1}} + \dots$$

$\mathcal{N} \rightarrow u^{\text{inc}}$



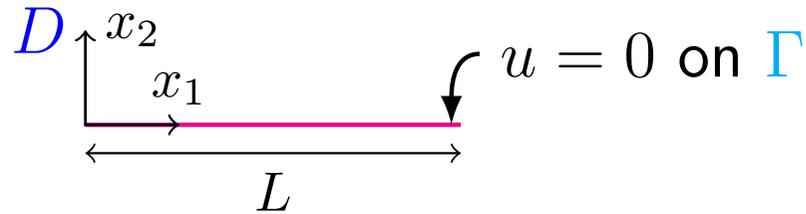
We will solve the BIE by a Galerkin BEM representing  $\left[\frac{\partial u}{\partial n}\right]$  as a sum of products of **known oscillatory functions** and **unknown non-oscillatory functions**.

Step A. Work out what this representation for  $\left[\frac{\partial u}{\partial n}\right]$  is - use **GTD**.

Above the screen ... and adding in the GTD terms ...

$$\frac{\partial u}{\partial n}(x) \approx 2 \frac{\partial u^{\text{inc}}}{\partial n}(x) + c_+ \frac{e^{ikx_1}}{\sqrt{kx_1}} + c_- \frac{e^{-ikx_1}}{\sqrt{k(L-x_1)}}$$

$\mathcal{N} \rightarrow u^{\text{inc}}$



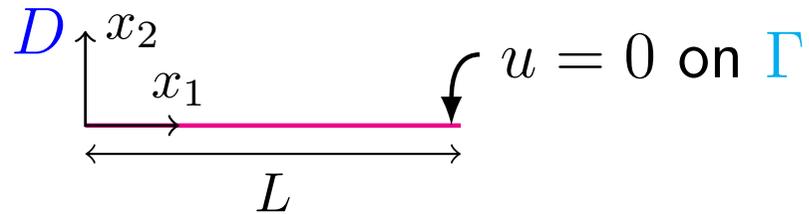
We will solve the BIE by a Galerkin BEM representing  $\left[\frac{\partial u}{\partial n}\right]$  as a sum of products of **known oscillatory functions** and **unknown non-oscillatory functions**.

Step A. Work out what this representation for  $\left[\frac{\partial u}{\partial n}\right]$  is - use **GTD**.

While below the screen

$$\frac{\partial u}{\partial n}(x) \approx c'_+ \frac{e^{ikx_1}}{\sqrt{kx_1}} + c'_- \frac{e^{-ikx_1}}{\sqrt{k(L-x_1)}}$$

$\mathcal{N} \rightarrow u^{\text{inc}}$



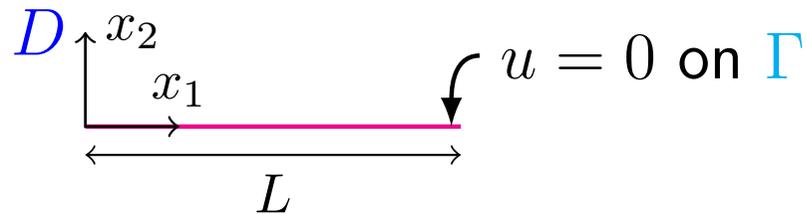
We will solve the BIE by a Galerkin BEM representing  $\left[\frac{\partial u}{\partial n}\right]$  as a sum of products of **known oscillatory functions** and **unknown non-oscillatory functions**.

Step A. Work out what this representation for  $\left[\frac{\partial u}{\partial n}\right]$  is - use **GTD**.

So

$$\left[\frac{\partial u}{\partial n}\right](x) \approx 2\frac{\partial u^{\text{inc}}}{\partial n}(x) + C_+ \frac{e^{ikx_1}}{\sqrt{kx_1}} + C_- \frac{e^{-ikx_1}}{\sqrt{k(L-x_1)}}.$$

$\mathcal{N} \rightarrow u^{\text{inc}}$



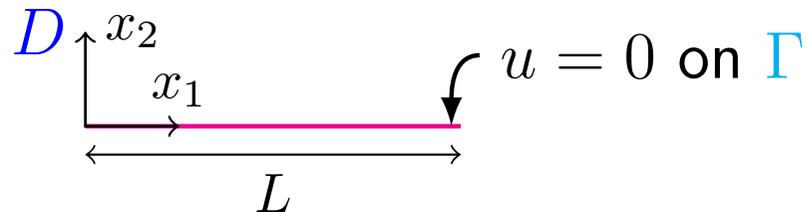
We will solve the BIE by a Galerkin BEM representing  $\left[\frac{\partial u}{\partial n}\right]$  as a sum of products of **known oscillatory functions** and **unknown non-oscillatory functions**.

Step A. Work out what this representation for  $\left[\frac{\partial u}{\partial n}\right]$  is - use **GTD**.

**... and our representation is ...**

$$\left[\frac{\partial u}{\partial n}\right](x) = 2\frac{\partial u^{\text{inc}}}{\partial n}(x) + e^{ikx_1} \underbrace{F_+(x_1)}_{\text{unknown}} + e^{-ikx_1} \underbrace{F_-(L-x_1)}_{\text{unknown}}$$

$\mathcal{N} \rightarrow u^{\text{inc}}$



We will represent  $\left[\frac{\partial u}{\partial n}\right]$  as a sum of products of **known oscillatory functions** and **unknown non-oscillatory functions**.

Step A. Work out what this representation for  $\left[\frac{\partial u}{\partial n}\right]$  is - use **GTD**.

**... and our representation is ...**

$$\left[\frac{\partial u}{\partial n}\right](x) = 2 \frac{\partial u^{\text{inc}}}{\partial n}(x) + e^{ikx_1} \underbrace{F_+(x_1)}_{\text{unknown}} + e^{-ikx_1} \underbrace{F_-(L-x_1)}_{\text{unknown}}$$

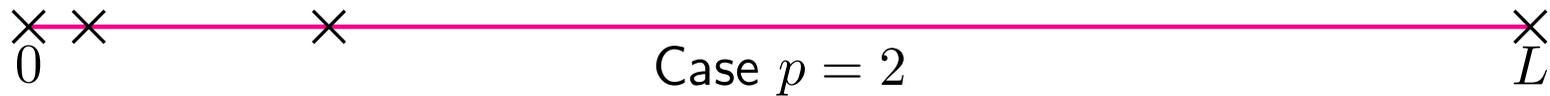
**Step B.** Show that  $F_{\pm}$  are **non-oscillatory** and choose **piecewise polynomial** approximation spaces for them.

**STEP B. RECAP – A PIECEWISE POLYNOMIAL QUIZ**

**How best to approximate  $F(t) = t^{1/2}$  on  $[0, L]$  with smallest  $L^\infty(0, L)$  error using piecewise polynomials of degree  $p$ ?**

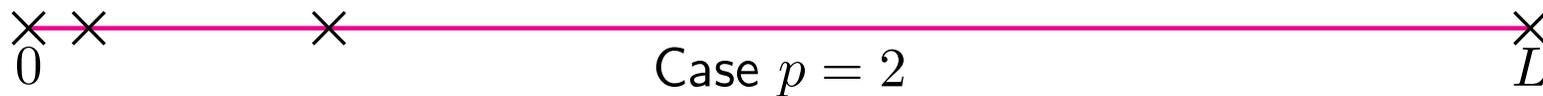


**How best to approximate  $F(t) = t^{1/2}$  on  $[0, L]$  with smallest  $L^\infty(0, L)$  error using piecewise polynomials of degree  $p$ ?**



**ANSWER! Grid-points  $0$  and  $L\alpha^j$ ,  $j = 0, \dots, p$ , with  $\alpha \approx 0.2$ , and polynomial of degree  $p$  on each subinterval.**

How best to approximate  $F(t) = t^{1/2}$  on  $[0, L]$  with smallest  $L^\infty(0, L)$  error using piecewise polynomials of degree  $p$ ?



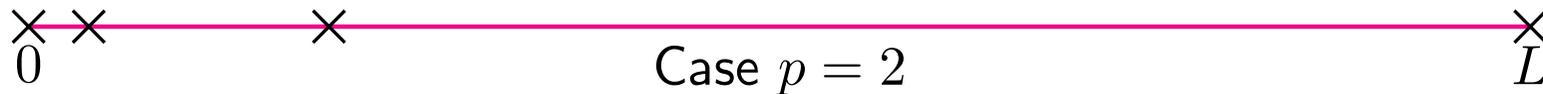
**ANSWER!** Grid-points  $0$  and  $L\alpha^j$ ,  $j = 0, \dots, p$ , with  $\alpha \approx 0.2$ , and polynomial of degree  $p$  on each subinterval.

This is standard *hp*-approximation on a geometrically graded mesh.

**How best to approximate  $F(t) = t^{-1/2}$  on  $[0, L]$  with smallest  $L^q(0, L)$  error ( $1 \leq q < 2$ ) using piecewise polynomials of degree  $p$ ?**



**How best to approximate  $F(t) = t^{-1/2}$  on  $[0, L]$  with smallest  $L^q(0, L)$  error ( $1 \leq q < 2$ ) using piecewise polynomials of degree  $p$ ?**



**ANSWER! Grid-points  $0$  and  $L\alpha^j$ ,  $j = 0, \dots, p$ , with  $\alpha \approx 0.2$ , and polynomial of degree  $p$  on each subinterval.**

This is standard *hp*-approximation on a geometrically graded mesh and

**minimum error  $\leq C \exp(-cp) = C \exp(-c\sqrt{N})$ ,**

**where  $N = D.O.F.$**

**Suppose that  $F(z)$  is analytic in  $\Re z > 0$  and**

$$|F(z)| \leq |z|^{-1/2}, \quad \Re z > 0.$$

**How best to approximate  $F$  on  $[0, L]$  with smallest  $L^q(0, L)$  error ( $1 \leq q < 2$ ) using piecewise polynomials of degree  $p$ ?**



Suppose that  $F(z)$  is analytic in  $\Re z > 0$  and

$$|F(z)| \leq |z|^{-1/2}, \quad \Re z > 0.$$

How best to approximate  $F$  on  $[0, L]$  with smallest  $L^q(0, L)$  error ( $1 \leq q < 2$ ) using piecewise polynomials of degree  $p$ ?



**ANSWER! Grid-points  $0$  and  $L\alpha^j$ ,  $j = 0, \dots, p$ , with  $\alpha \approx 0.2$ , and polynomial of degree  $p$  on each subinterval.**

This is standard *hp*-approximation on a geometrically graded mesh and

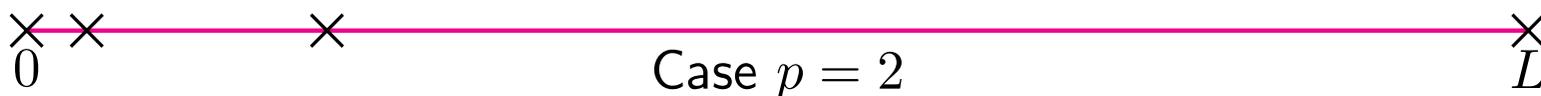
$$\text{minimum error} \leq C \exp(-cp) = C \exp(-c\sqrt{N}),$$

where  $N = D.O.F.$ .

Suppose that  $F(z)$  is analytic in  $\Re z > 0$  and

$$|F(z)| \leq |z|^{-1/2}, \quad \Re z > 0.$$

How best to approximate  $F$  on  $[0, L]$  with smallest  $\tilde{H}^{-1/2}(\Gamma)$  error using piecewise polynomials of degree  $p$ ?



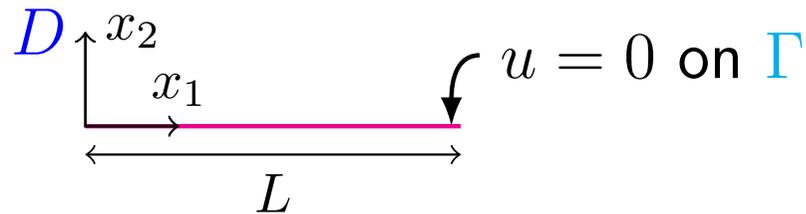
**ANSWER! Grid-points  $0$  and  $L\alpha^j$ ,  $j = 0, \dots, p$ , with  $\alpha \approx 0.2$ , and polynomial of degree  $p$  on each subinterval.**

**This follows since  $\tilde{H}^{-1/2}(\Gamma)$  is continuously embedded in  $L^q(\Gamma) = L^q(0, L)$  for  $q > 1$ , and hence**

**minimum error  $\leq C \exp(-cp) = C \exp(-c\sqrt{N})$ ,**

**where  $N = D.O.F.$**

$\mathcal{N}_{\rightarrow} u^{\text{inc}}$



**Non-Oscillatory Theorem.** For some  $F_{\pm}(z)$ , analytic in  $\Re z > 0$  with

$$|F_{\pm}(z)| \leq C k^{3/2} |z|^{-1/2}, \quad \Re z > 0,$$

it holds that

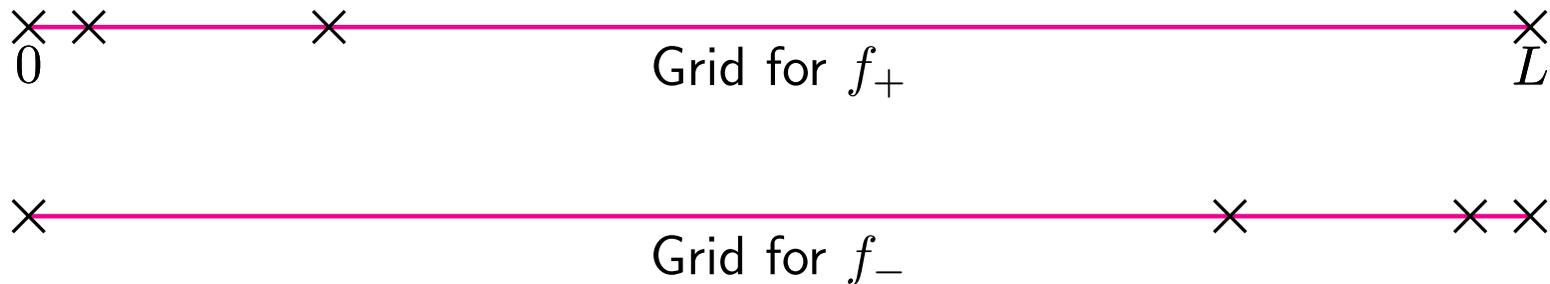
$$\left[ \frac{\partial u}{\partial n} \right] (x) = 2 \frac{\partial u^{\text{inc}}}{\partial n} (x) + e^{ikx_1} F_+(x_1) + e^{-ikx_1} F_-(L - x_1), \quad x \in \Gamma.$$

*Proof.* Hewett, Langdon, C-W, to appear *IMA J. Numer. Anal.*

For  $x \in \Gamma$ ,

$$\begin{aligned} \left[ \frac{\partial u}{\partial n} \right] (x) &= 2 \frac{\partial u^{\text{inc}}}{\partial n} (x) + e^{ikx_1} F_+(x_1) + e^{-ikx_1} F_-(L - x_1) \\ &\approx 2 \frac{\partial u^{\text{inc}}}{\partial n} (x) + e^{ikx_1} f_+(x_1) + e^{-ikx_1} f_-(x_1), \end{aligned}$$

where  $f_+$  and  $f_-$  are piecewise polynomials of degree  $p$  on geometrically graded meshes, each with  $p$  intervals: i.e.,  $hp$ -approximation.

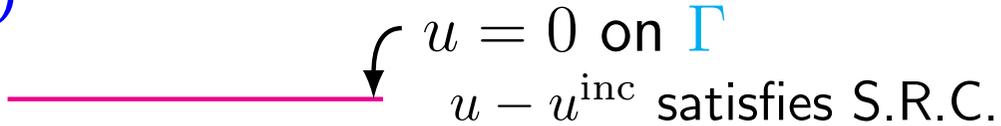


**This is our Galerkin approximation space.**

## A Case Study for Numerical-Asymptotic Methods

$$\mathcal{W} \rightarrow u^{\text{inc}} \quad \Delta u + k^2 u = 0$$

$D$



$u = 0$  on  $\Gamma$   
 $u - u^{\text{inc}}$  satisfies S.R.C.

We will solve

$$0 = u^{\text{inc}}(x) + \int_{\Gamma} \left[ \frac{\partial u}{\partial n} \right] (y) \Phi(x, y) ds(y), \quad x \in \Gamma,$$

in operator notation

$$S_k \left[ \frac{\partial u}{\partial n} \right] = -u^{\text{inc}}|_{\Gamma}, \quad \text{(BIE)}$$

by a Galerkin BEM using this ansatz.

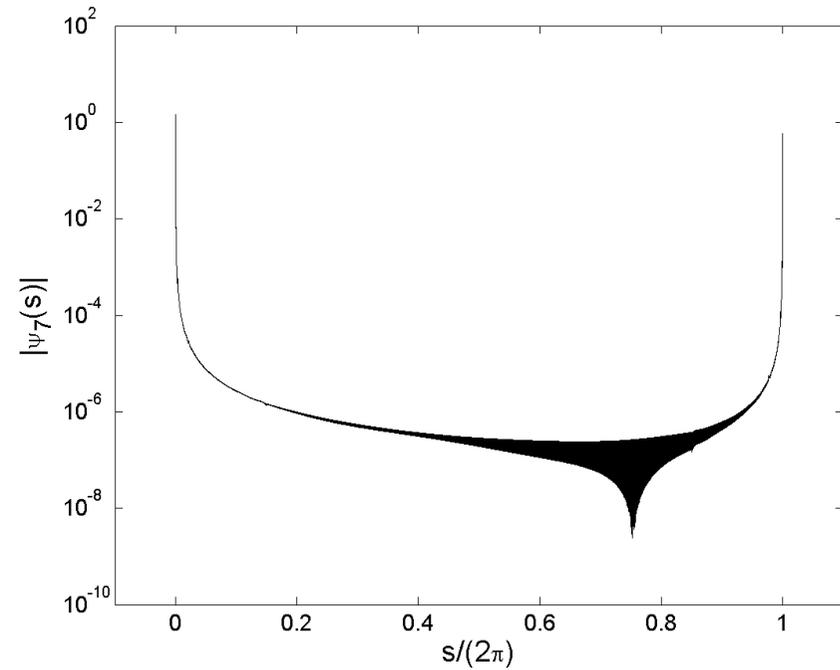
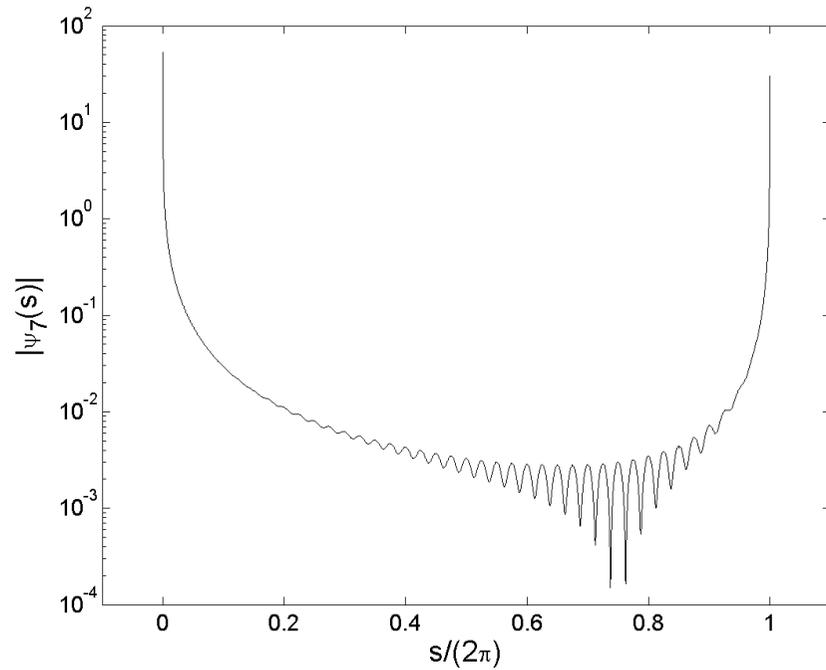
**STEP C. IMPLEMENT IT AND SEE THAT THE COST  
IS  $O(1)$  AS  $k \rightarrow \infty$ !**

Difficulty (and main cost) is assembly of the matrix  $[a_{mn}]$  which requires **2D highly oscillatory integrals**:

$$a_{mn} = \int_{\Gamma_m} \int_{\Gamma_n} H_0^{(1)}(k|x_1 - y_1|) \exp(\pm i k x_1 \pm i k y_1) p_m(x_1) p_n(y_1) dx_1 dy_1,$$

where  $p_m$  and  $p_n$  are polynomials supported on elements  $\Gamma_m$  and  $\Gamma_n$ .

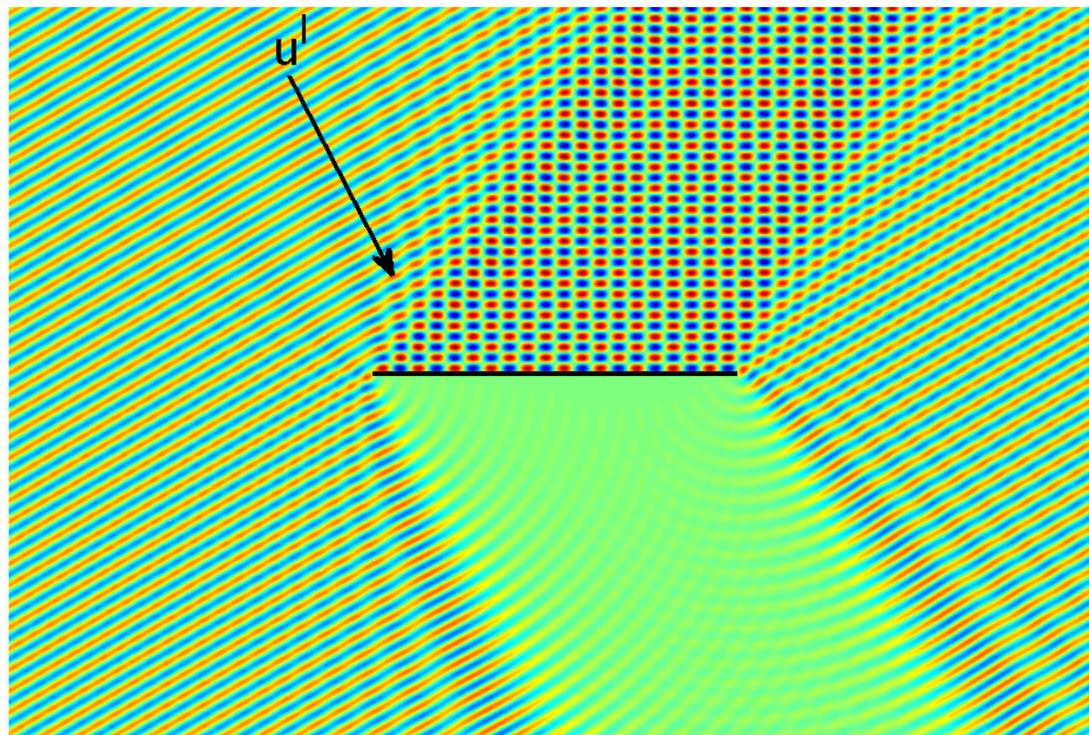
For details of our Filon quadrature see Hewett, Langdon, C-W (2014).



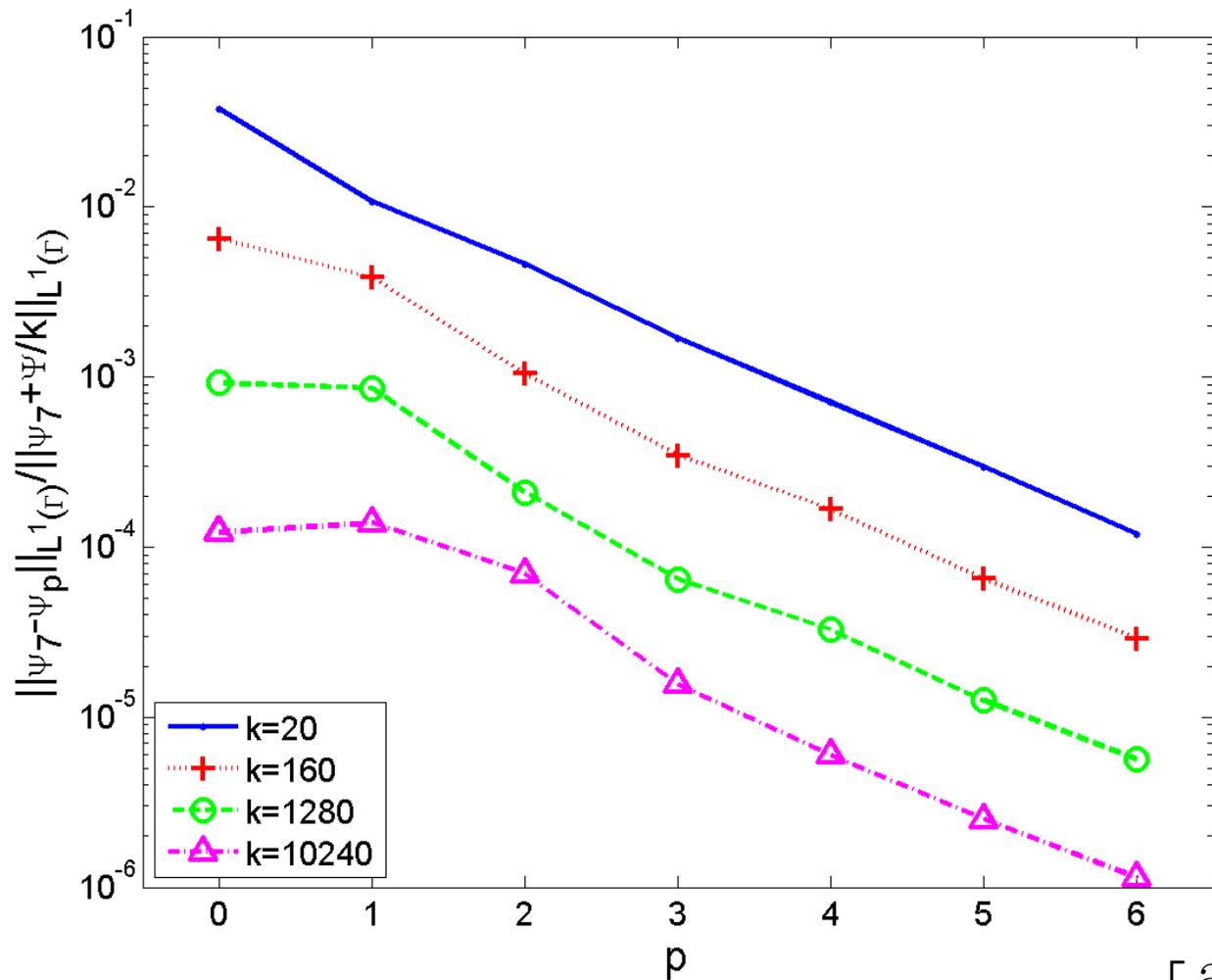
Plots of the amplitude of the diffracted component, i.e.

$$\left| \left[ \frac{\partial u}{\partial n} \right] (x) - 2 \frac{\partial u^{\text{inc}}}{\partial n} (x) \right| \quad \text{against} \quad \frac{x_1}{L},$$

for  $\frac{L}{\lambda} = 20$  (left),  $\frac{L}{\lambda} = 10240$  (right).



$\Re u$  in  $D$  for  $\frac{L}{\lambda} = 20$ .



Relative  $L^1(\Gamma)$  errors in Galerkin approximation to  $\left[ \frac{\partial u}{\partial n} \right]$ .

$\frac{L}{\lambda}$	$\lambda N/L$	$\frac{\ [\partial u/\partial n] - \phi_{64}\ _{L^1(\Gamma)}}{\ [\partial u/\partial n]\ _{L^1(\Gamma)}}$	cpu time (secs)
10	$6.40 \times 10^0$	$1.38 \times 10^{-2}$	47
40	$1.60 \times 10^0$	$1.40 \times 10^{-2}$	42
160	$4.00 \times 10^{-1}$	$1.40 \times 10^{-2}$	47
640	$1.00 \times 10^{-1}$	$1.39 \times 10^{-2}$	42
2560	$2.50 \times 10^{-2}$	$1.38 \times 10^{-2}$	42
10240	$6.25 \times 10^{-3}$	$1.37 \times 10^{-2}$	40

Relative  $L^1(\Gamma)$  error in computing  $[\partial u/\partial n]$ :  
64 degrees of freedom, grazing incidence.

**STEP D. TRY TO PROVE THAT THE METHOD IS  $O(1)$  BY THEOREMS ABOUT THE  $k$ -DEPENDENCE OF EVERYTHING!**

## Error Analysis

We are solving by a Galerkin BEM

$$0 = u^{\text{inc}}(x) + \int_{\Gamma} \left[ \frac{\partial u}{\partial n} \right] (y) \Phi(x, y) ds(y), \quad x \in \Gamma,$$

in operator notation

$$S_{\mathbf{k}} \left[ \frac{\partial u}{\partial n} \right] = -u^{\text{inc}}|_{\Gamma}.$$

By explicit Fourier analysis we can show, for a general planar screen (in 2D or 3D),  **$\mathbf{k}$ -explicit coercivity** and **continuity** of  $S_{\mathbf{k}}$ : for some  $C, \alpha > 0$  independent of  $\mathbf{k}$  and  $L$ ,

$$\|S_{\mathbf{k}}\| \leq C(\mathbf{k}L)^{1/2}, \quad |\langle S_{\mathbf{k}}\phi, \phi \rangle| \geq \alpha \|\phi\|_{\tilde{H}^{-1/2}(\Gamma)}^2,$$

where  $\|S_{\mathbf{k}}\|$  is the norm of  $S_{\mathbf{k}} : \tilde{H}^{-1/2}(\Gamma) \rightarrow H^{1/2}(\Gamma)$  and  $\langle \cdot, \cdot \rangle$  is the  $L^2(\Gamma)$  inner product.

## Error Analysis

We are solving by a Galerkin BEM

$$0 = u^{\text{inc}}(x) + \int_{\Gamma} \left[ \frac{\partial u}{\partial n} \right] (y) \Phi(x, y) ds(y), \quad x \in \Gamma,$$

in operator notation

$$S_{\mathbf{k}} \left[ \frac{\partial u}{\partial n} \right] = -u^{\text{inc}}|_{\Gamma}.$$

By explicit Fourier analysis we can show, for a general planar screen (in 2D or 3D),  **$\mathbf{k}$ -explicit coercivity** and **continuity** of  $S_{\mathbf{k}}$ : for some  $C, \alpha > 0$  independent of  $\mathbf{k}$  and  $L$ ,

$$\|S_{\mathbf{k}}\| \leq C(\mathbf{k}L)^{1/2}, \quad |\langle S_{\mathbf{k}}\phi, \phi \rangle| \geq \alpha \|\phi\|_{\tilde{H}^{-1/2}(\Gamma)}^2,$$

where  $\|S_{\mathbf{k}}\|$  is the norm of  $S_{\mathbf{k}} : \tilde{H}^{-1/2}(\Gamma) \rightarrow H^{1/2}(\Gamma)$  and  $\langle \cdot, \cdot \rangle$  is the  $L^2(\Gamma)$  inner product. **Surprisingly definite for Helmholtz!**

## Error Analysis

$$S_{\mathbf{k}} \begin{bmatrix} \frac{\partial u}{\partial n} \end{bmatrix} = -u^{\text{inc}}|_{\Gamma}.$$

By explicit Fourier analysis we can show, for a general planar screen (in 2D or 3D),  $\mathbf{k}$ -explicit coercivity and continuity of  $S_{\mathbf{k}}$ : for some  $C, \alpha > 0$  independent of  $\mathbf{k}$  and  $L$ ,

$$\|S_{\mathbf{k}}\| \leq C(\mathbf{k}L)^{1/2}, \quad |\langle S_{\mathbf{k}}\phi, \phi \rangle| \geq \alpha \|\phi\|_{\tilde{H}^{-1/2}(\Gamma)}^2,$$

where  $\|S_{\mathbf{k}}\|$  is the norm of  $S_{\mathbf{k}} : \tilde{H}^{-1/2}(\Gamma) \rightarrow H^{1/2}(\Gamma)$ .

By **Céa's lemma** the Galerkin solution  $\phi_N$  is well-defined and

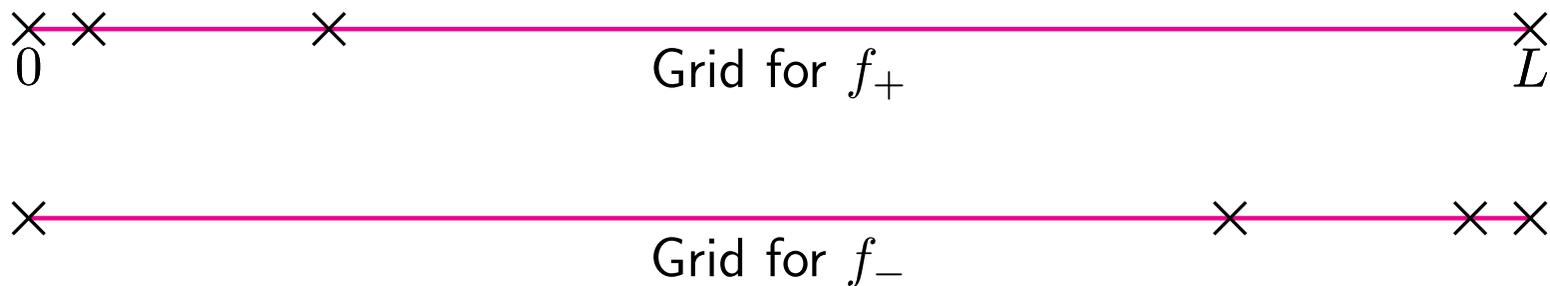
$$\left\| \begin{bmatrix} \frac{\partial u}{\partial n} \end{bmatrix} - \phi_N \right\|_{\tilde{H}^{-1/2}(\Gamma)} \leq \frac{C}{\alpha} (\mathbf{k}L)^{1/2} \inf_{\psi_N} \left\| \begin{bmatrix} \frac{\partial u}{\partial n} \end{bmatrix} - \psi_N \right\|_{\tilde{H}^{-1/2}(\Gamma)},$$

where the infimum is taken over all  $\psi_N$  in the  $N$ -dimensional Galerkin subspace.

For  $x \in \Gamma$ ,

$$\begin{aligned} \left[ \frac{\partial u}{\partial n} \right] (x) &= 2 \frac{\partial u^{\text{inc}}}{\partial n} (x) + e^{ikx_1} F_+(x_1) + e^{-ikx_1} F_-(L - x_1) \\ &\approx 2 \frac{\partial u^{\text{inc}}}{\partial n} (x) + e^{ikx_1} f_+(x_1) + e^{-ikx_1} f_-(x_1), \end{aligned}$$

where  $f_+$  and  $f_-$  are piecewise polynomials of degree  $p$  on geometrically graded meshes, each with  $p$  intervals: i.e.,  $hp$ -approximation.



**This is our Galerkin approximation space.**

$$\left[ \frac{\partial u}{\partial n} \right] (x) \approx 2 \frac{\partial u^{\text{inc}}}{\partial n} (x) + e^{ikx_1} f_+(x_1) + e^{-ikx_1} f_-(x_1),$$

where  $f_+$  and  $f_-$  are piecewise polynomials of degree  $p$  on geometrically graded meshes, each with  $p$  intervals: i.e.,  $hp$ -approximation.

**Theorem** If  $\phi_N$  is the best  $\tilde{H}^{-1/2}(\Gamma)$  approximation to  $\left[ \frac{\partial u}{\partial n} \right]$  of this form, then

$$\left\| \left[ \frac{\partial u}{\partial n} \right] - \phi_N \right\|_{\tilde{H}^{-1/2}(\Gamma)} \leq C k^{3/2} (\log k)^{1/2} \exp(-c\sqrt{N}),$$

where  $C$  and  $c$  depend (only) on  $\Gamma$ , and  $N \propto p^2$  is the number of D.O.F.

$$\left[ \frac{\partial u}{\partial n} \right] (x) \approx 2 \frac{\partial u^{\text{inc}}}{\partial n} (x) + e^{ikx_1} f_+(x_1) + e^{-ikx_1} f_-(x_1),$$

where  $f_+$  and  $f_-$  are piecewise polynomials of degree  $p$  on geometrically graded meshes, each with  $p$  intervals: i.e.,  $hp$ -approximation.

**Theorem** If  $\phi_N$  is the Galerkin approximation to  $\left[ \frac{\partial u}{\partial n} \right]$  of this form, then

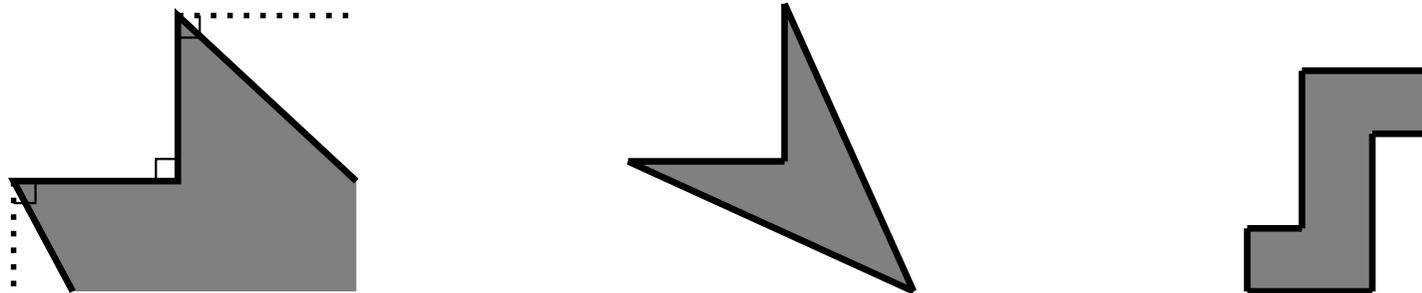
$$\left\| \left[ \frac{\partial u}{\partial n} \right] - \phi_N \right\|_{\tilde{H}^{-1/2}(\Gamma)} \leq C k^2 (\log k)^{1/2} \exp(-c\sqrt{N}),$$

where  $C$  and  $c$  depend (only) on  $\Gamma$ , and  $N \propto p^2$  is the number of D.O.F.

### **3. EXTENSIONS TO OTHER GEOMETRIES AND 3D**

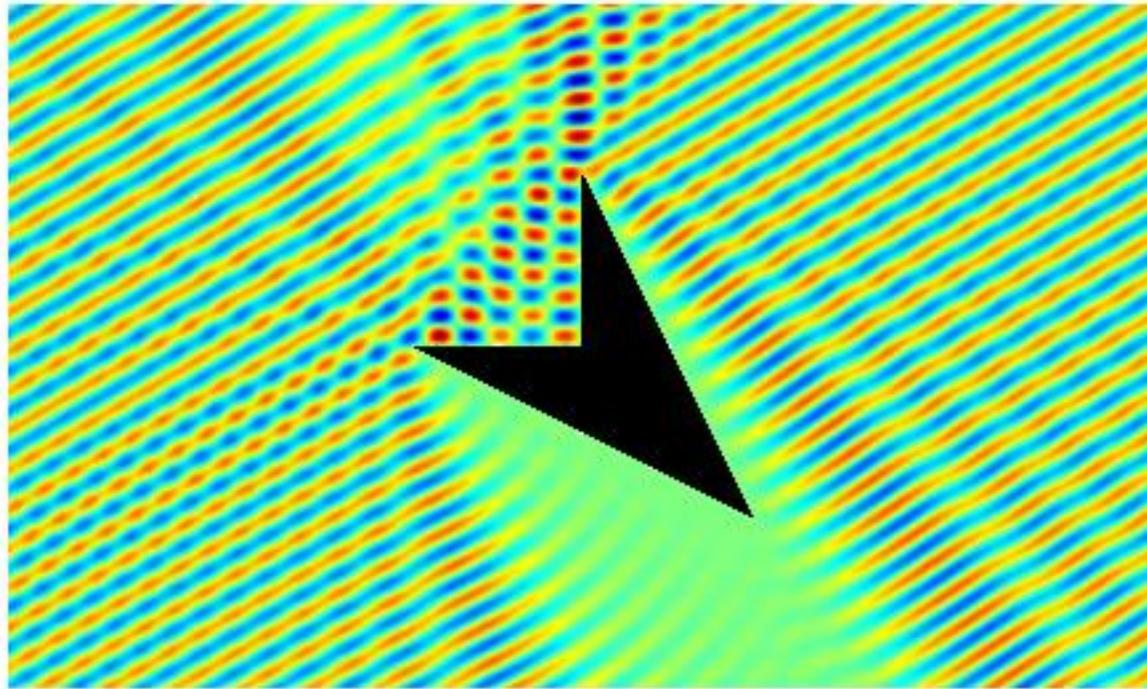
## Polygons: Convex and Non-Convex

C-W, Hewett, Langdon, Twigger, *Numer Math* (2014)

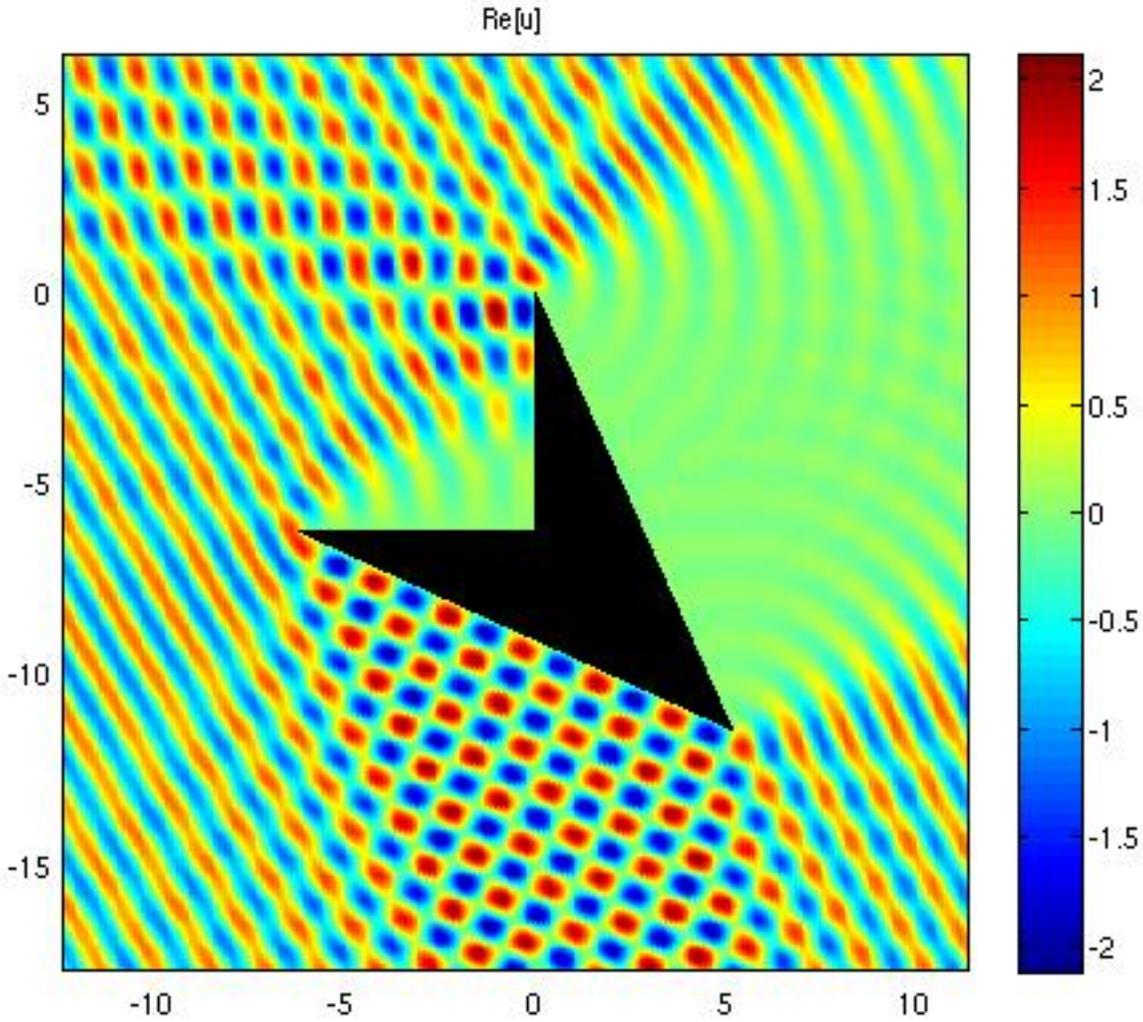


We can, using GO/GTD, design an approximation space for  $\frac{\partial u}{\partial n}$  which provably needs only  $O(\log^2 k)$  degrees of freedom as  $k \rightarrow \infty$  and in experiments only  $O(1)$ .

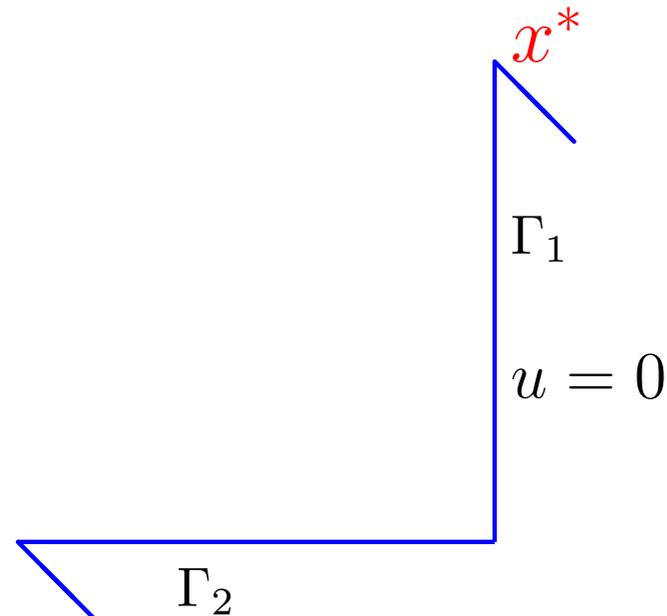
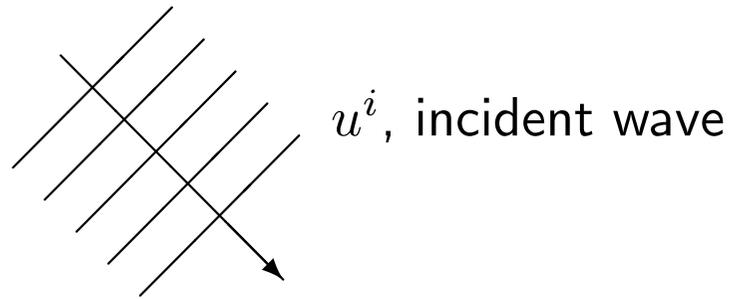
## Solution Behaviour: $\Re u$



# Solution Behaviour: $\Re u$



## Solution Behaviour on $\Gamma_2$



On  $\Gamma_2$ ,

$$\frac{\partial u}{\partial n} = \text{known} + e^{ik|x-x^*|} F(x_1) + e^{ikx_1} F_+(x_1) + e^{-ikx_1} F_-(x_1)$$

where 'known' = Fresnel integral and  $F$  is **analytic** and bounded in fixed neighbourhood of  $\Gamma_2$ , and again  $N = O(\log^2 k)$  as  $k \rightarrow \infty$  is provably enough.

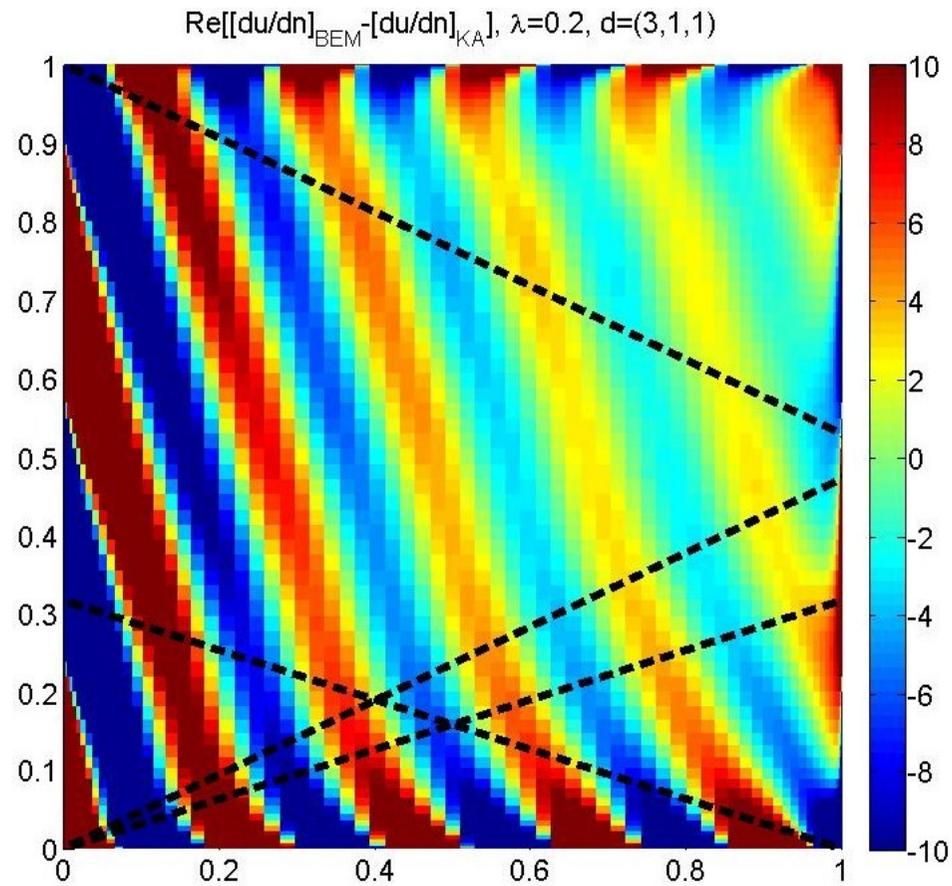
*hp*-BEM Based on this Ansatz

$k$	dof	dof per $\lambda$	$L^2$ error	Relative $L^2$ error
5	320	10.7	2.09e-2	1.51e-2
10	320	5.3	1.07e-2	1.11e-2
20	320	2.7	4.60e-3	6.91e-3
40	320	1.3	3.13e-3	6.83e-3

C-W, Langdon, Hewett, Twigger, *Numer Math* (2014).

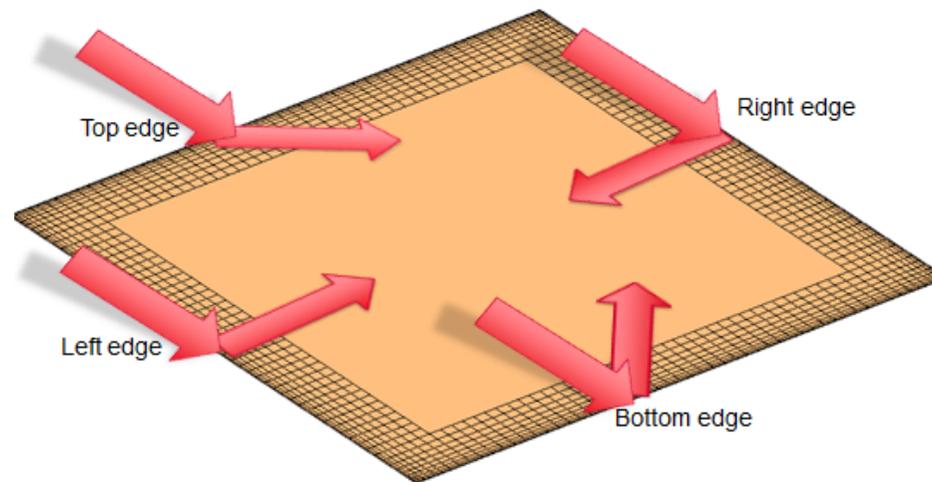
### 3D Thin Screen: Square Plate

Hargreaves, Hewett, Langdon, Lam,  
EPSRC project Reading/Salford

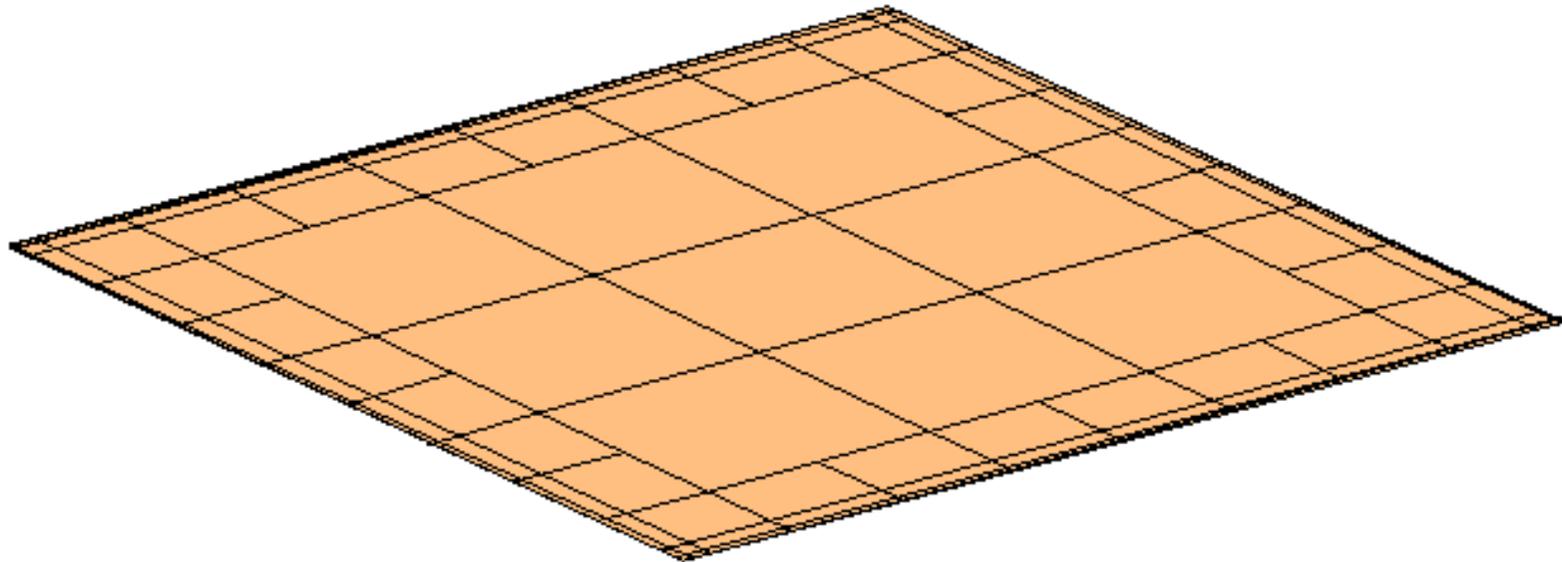


## Approximation Methodology

- Subtract leading order oscillatory behaviour (incident field).
- Small conventional elements around the rim (to represent singular behaviour at edge).
- Large hybrid elements in the centre; basis functions are plane waves multiplied by polynomial basis functions (order  $p$ ).
- Phase functions on hybrid elements correspond to first order diffraction directions (“edge plane waves”).

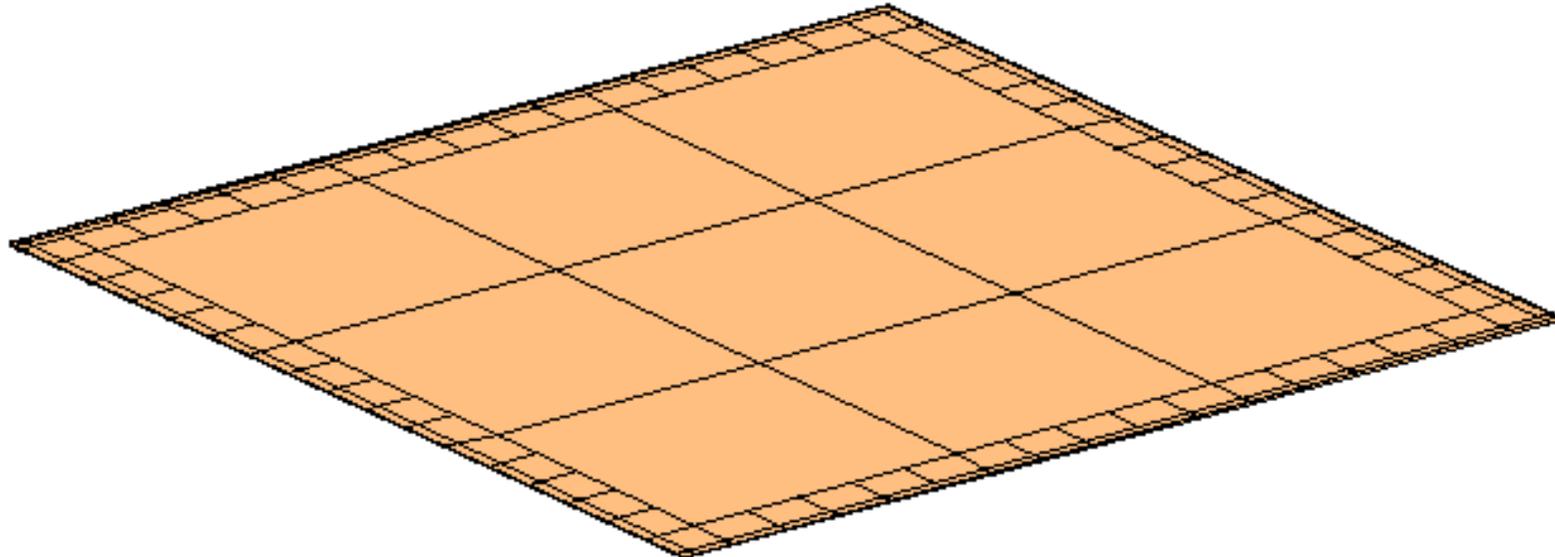


Mesh and required number of DOFS,  $k = 5$



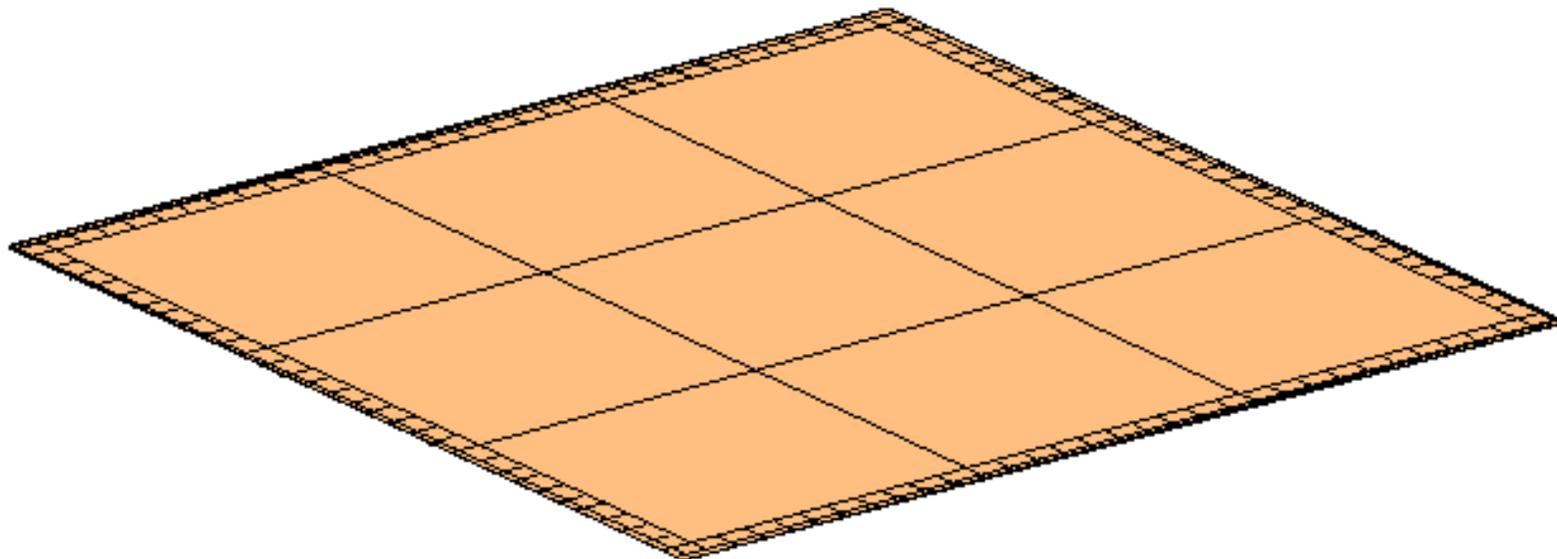
#DOF	Constant	Linear	Quadratic	Cubic
Regular	196	576	1,296	2,500
Hybrid	88	396	1,044	2,176

Mesh and required number of DOFS,  $k = 10$



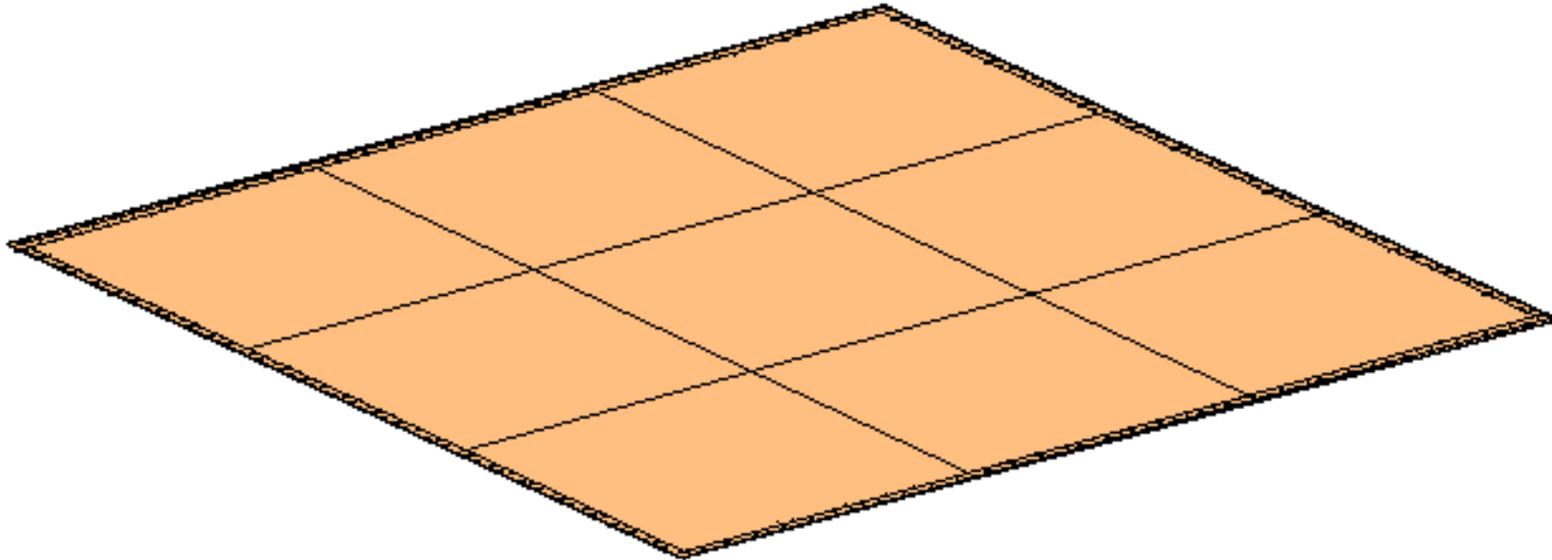
#DOF	Constant	Linear	Quadratic	Cubic
Regular	1,024	2,601	5,184	9,025
Hybrid	160	720	1,908	3,976

Mesh and required number of DOFS,  $k = 20$



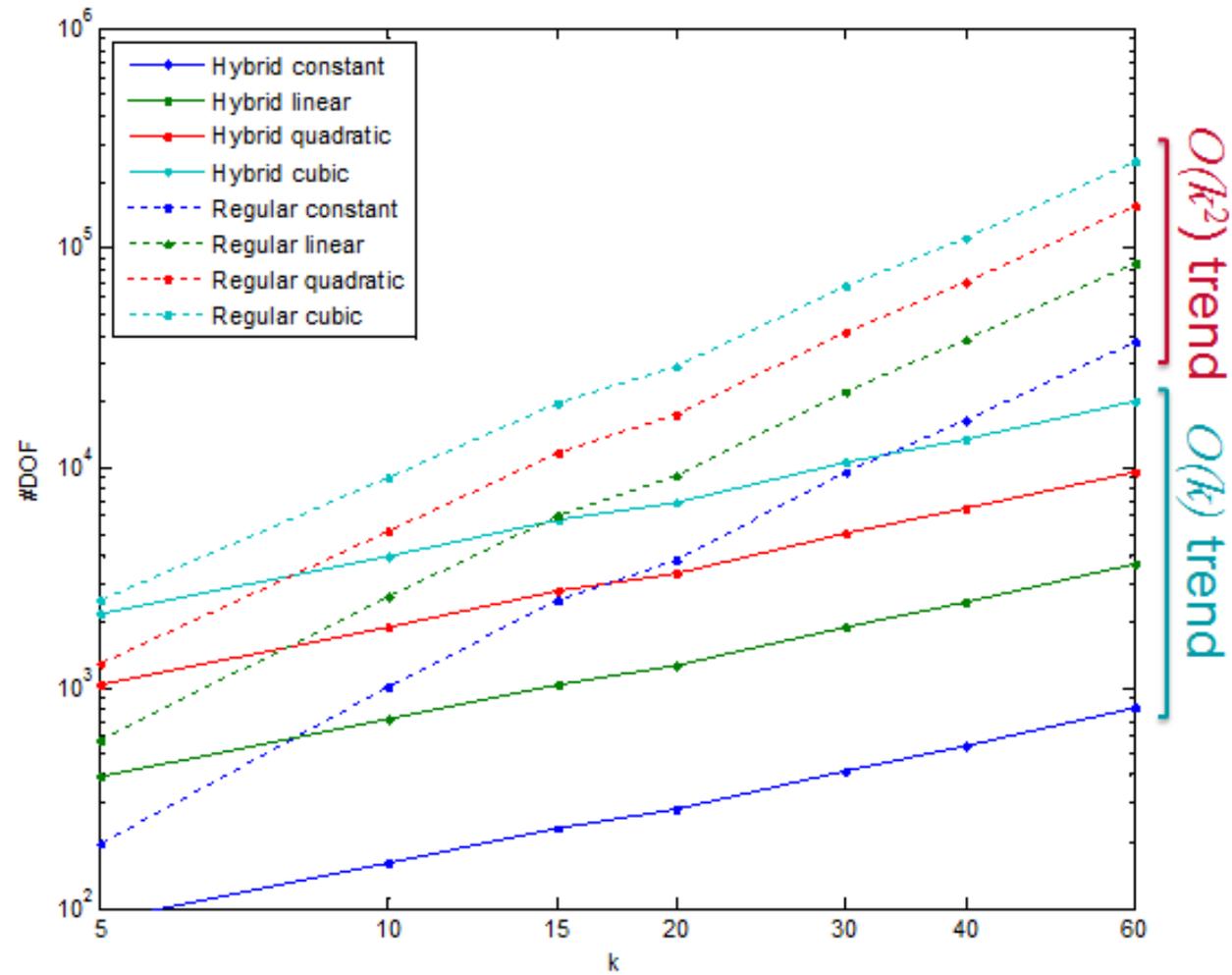
#DOF	Constant	Linear	Quadratic	Cubic
Regular	3,844	9,216	17,424	28,900
Hybrid	280	1,260	3,348	6,976

Mesh and required number of DOFS,  $k = 40$

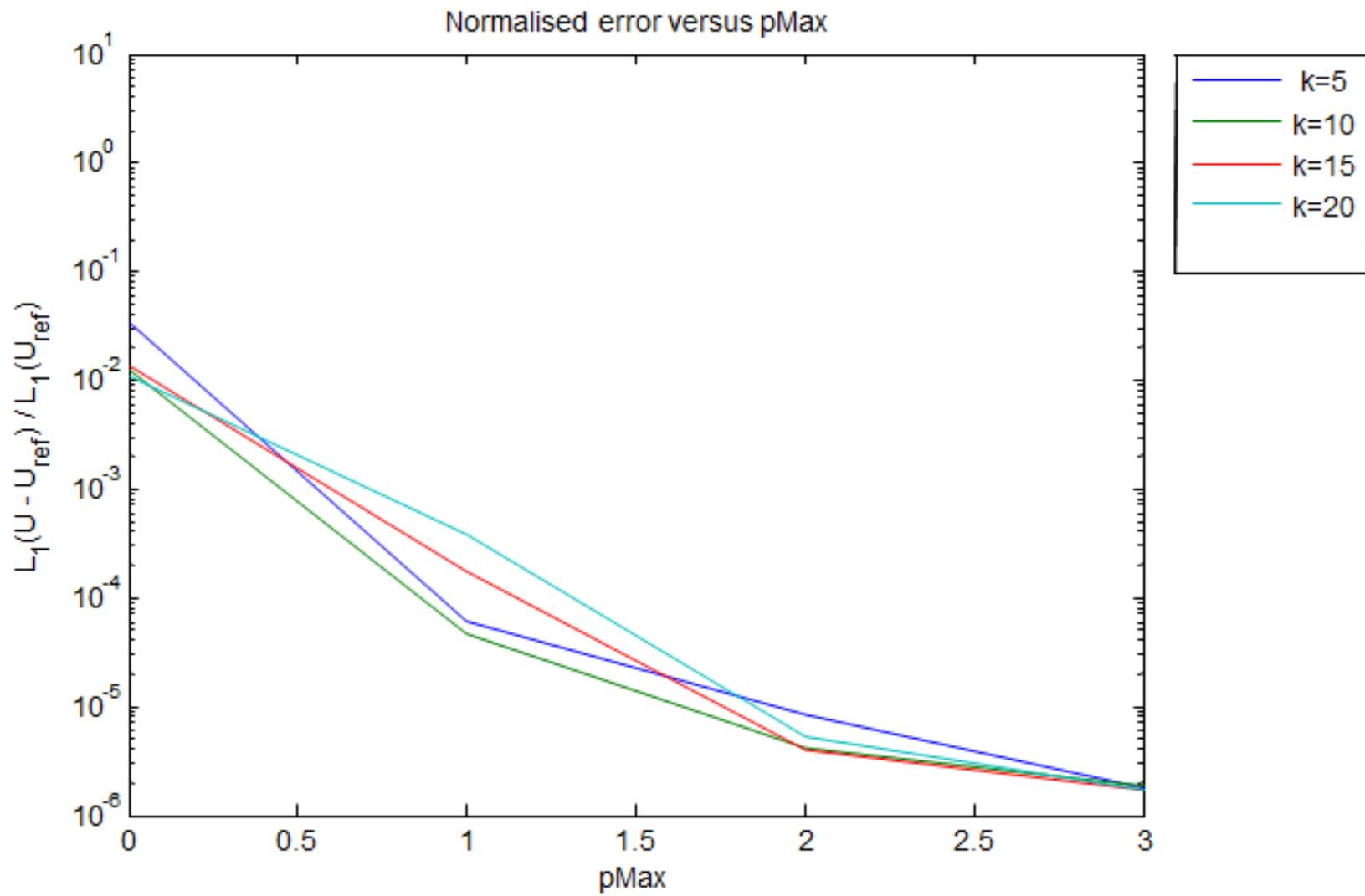


#DOF	Constant	Linear	Quadratic	Cubic
Regular	16,384	38,025	69,696	112,225
Hybrid	544	2,448	6,516	13,576

## Degrees of freedom trend



# Convergence of hybrid scheme



## Other Geometries

- **Smooth convex obstacles:** see Bruno et al. *Phil. Trans R. Soc.* (2004), Dominguez, Graham et al *Numer. Math.* (2007), Huybrechs & Vandewalle *SISC* (2007)
- **Piecewise smooth convex polygons:** see Langdon, Mokgolele, C-W *J. Comp. Appl. Math* (2010)
- **Inhomogeneous impedance plane: outdoor noise propagation:** see C-W, Langdon *Phil. Trans R. Soc.* (2004), Langdon & C-W *SINUM* (2006)
- **Penetrable scatterers:** see Groth, Hewett, Langdon *IMA J. Appl. Math.* (2014)

## Recap

1. Green's Representation Theorem and **boundary integral equations** (for Helmholtz)

2. A Case study for **numerical-asymptotic** methods: **the thin screen**

Step A. Represent the unknown as sum of products of **known oscillatory** and **unknown non-oscillatory** functions using **GO/GTD**.

Step B. Decide on the approximation space - combine **HF asymptotics** with *hp*-**approximation theory**

Step C. Implement it and see that (we hope) the cost is  $O(1)$  as  $k \rightarrow \infty$ !

Step D. Try to prove this by theorems about the  $k$ -dependence of everything!

3. **Other geometries and 3D**

## References

### Two review papers:

- “Numerical-asymptotic boundary integral methods in high-frequency acoustic scattering”, C-W, I.G. Graham, S Langdon, & E.A. Spence, *Acta Numerica* (2012).
- **“Acoustic scattering: high frequency boundary element methods and unified transform methods”**, C-W & Langdon, to appear (preprint on Researchgate).

Note **Unified transform methods**  $\approx$  **WBM** from Daan’s talk.