



# Discrete Calculus Methods and their Application to Fluid Dynamics

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# Context

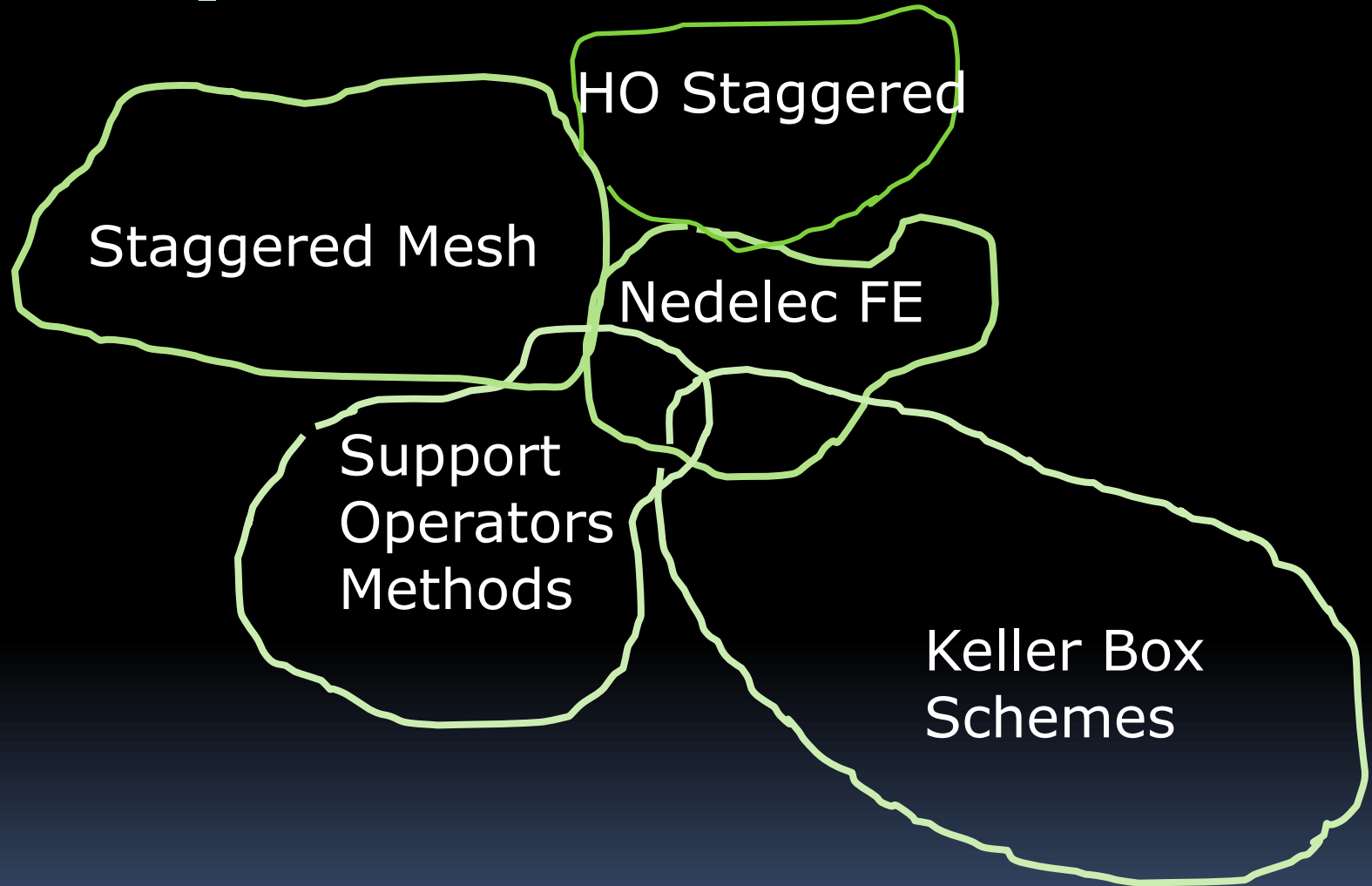
## Mathematical Tools

Order of Accuracy  
Stability  
Consistency

## Physical Requirements

Conservation  
(mass, momentum, total energy)  
Secondary Conservation  
(kinetic energy, vorticity, entropy)  
Unphysical Modes (pressure)  
Wave propagation (direction)  
Eigenmodes/Resonance

# Many Mimetic Methods



**Methods that capture physics well.**

# Why ?

**Why do some methods capture physics well?**

**They use exact discretization**

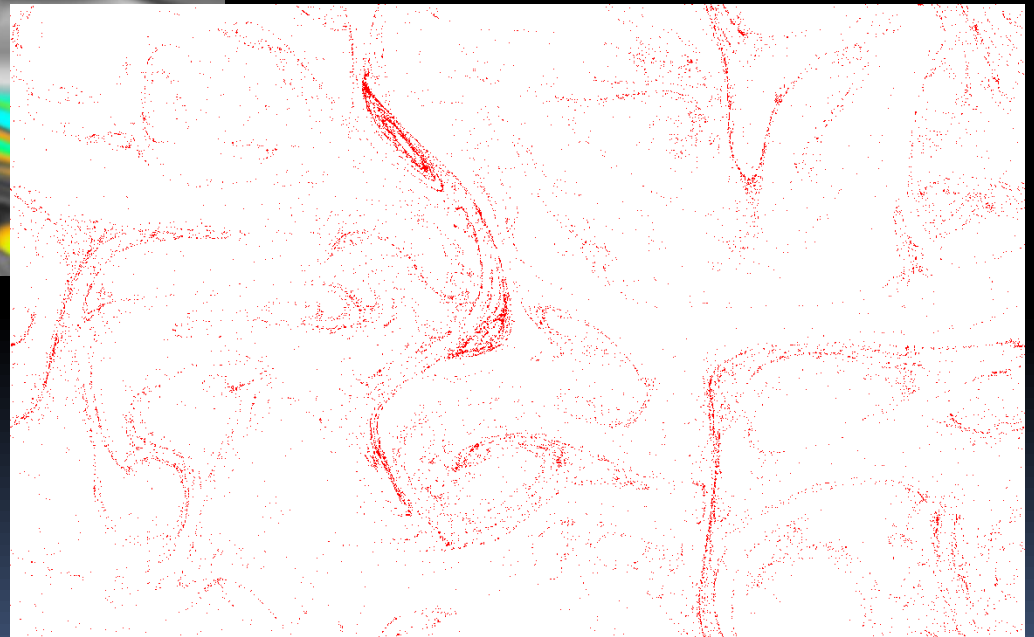
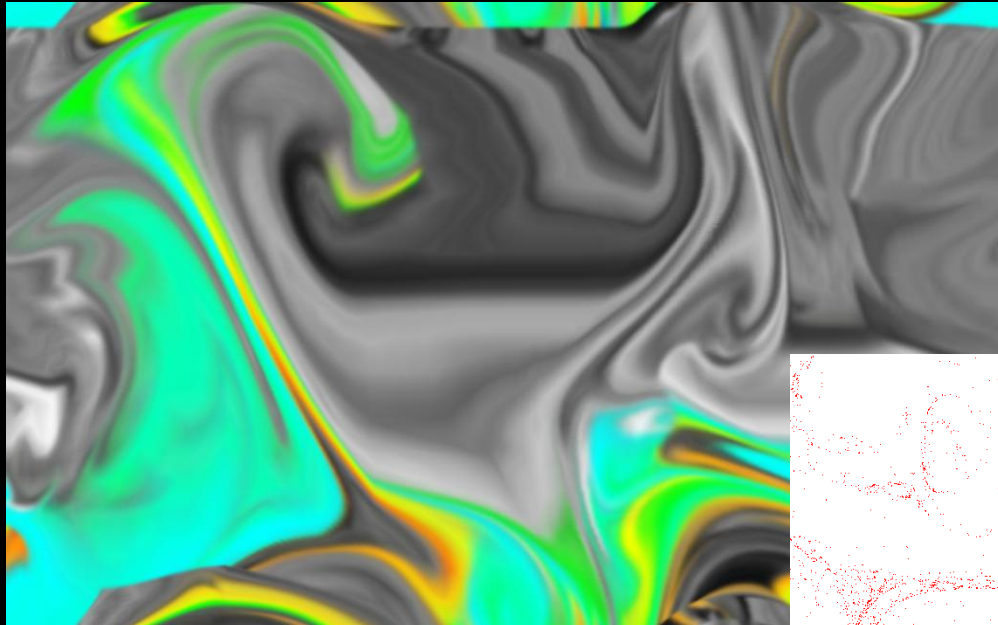
**What do they have in common?**

**All are Discrete Calculus Methods**

**Can you make a numerical method so that is mimetic from its design?**

**Yes (I think)**

# Discretization



**Take a continuous problem to a finite dimensional one.**

# Discrete Calculus: Part 1

## Exact Discretization

$$\frac{\partial a}{\partial t} + \nabla \cdot \mathbf{b} = 0 \quad \longrightarrow \quad \begin{bmatrix} \mathbf{A} & \mathbf{B} \end{bmatrix} \begin{pmatrix} \vec{\mathbf{a}} \\ \vec{\mathbf{b}} \end{pmatrix} = \vec{\mathbf{r}}$$

**Infinite Dimensional**

**Finite Dimensional**

**Partial Differential Eqn.**

**Matrix Problem**

**Basic unknowns are integral quantities.  
Collect infinite data into finite groups.**

# Discrete Calculus: Part 2

**Solution requires Approximation**

$$\begin{bmatrix} \mathbf{A} & \mathbf{B} \end{bmatrix} \begin{pmatrix} \vec{\mathbf{a}} \\ \vec{\mathbf{b}} \end{pmatrix} = \vec{\mathbf{r}} \quad \Rightarrow \quad \begin{bmatrix} \mathbf{A} & \mathbf{B} \\ \mathbf{C} & \mathbf{D} \end{bmatrix} \begin{pmatrix} \vec{\mathbf{a}} \\ \vec{\mathbf{b}} \end{pmatrix} = \begin{pmatrix} \vec{\mathbf{r}} \\ \mathbf{0} \end{pmatrix}$$

**Underdetermined**

**Unique Square**

**Relate discrete unknowns to each other.  
This relation is a material law.  
Also related to an dual mesh interpolation.  
Also related to inner products**

# Implications

**Exact Discretization means that:**

- **Calculus is exact.**
- **Physics is exact.**
- **Method is mimetic.**

**Discrete Solution Requires:**

- **Approximation of material laws.**
- **Interpolation between meshes.**
- **Assumptions about the solution.**



# Example: Heat Equation

$$\frac{\partial T}{\partial t} = \alpha \nabla^2 T$$

$$\frac{\partial(\rho CT)}{\partial t} = -\nabla \cdot \vec{\mathbf{q}}$$

Conservation of Energy

$$\vec{\mathbf{q}} = -k \nabla T$$

Fourier's Heat Flux Law

**Figure out what should NOT be approximated and what is already an approximation (Tonti).**

# Heat Eqn: Fully Separated

## Heat Equation

$$\frac{\partial i}{\partial t} = -\nabla \cdot \mathbf{q} \quad \text{Conservation of Energy (Physics)}$$

$$i \approx \rho C T \quad \text{Perfectly Caloric Material (Mat.)}$$

$$\mathbf{q} \approx -k \mathbf{g} \quad \text{Fourier's Heat Flux Law (Mat.)}$$

$$\mathbf{g} = \nabla T \quad \text{Def. of Gradient (Math)}$$

- **Discretize Physics and Math - exactly.**
- **Approximate Material laws - using interpolation.**

# Exact Discretization

$$I_{\tilde{c}}^{n+1} - I_{\tilde{c}}^n = - \sum_{\text{faces}} \bar{Q}_{\tilde{f}}^{\text{out}}$$

Gauss' Theorem (with time)  
Like FV: Not very original.

$$I_{\tilde{c}}^{n+1} = \int_{\Omega_{\tilde{c}}^{n+1}} idV \quad \bar{Q}_{\tilde{f}}^{\text{out}} = \int_{t^n}^{t^{n+1}} dt \int_{\partial\Omega_{\tilde{f}}} \mathbf{q} \cdot \mathbf{n}^{\text{out}} dA$$

$$\mathbf{g}_e^{n+1} = \int_{v_1}^{v_2} \mathbf{g} \cdot d\mathbf{l} \Big|^{n+1} = \int_{v_1}^{v_2} \nabla T \cdot d\mathbf{l} \Big|^{n+1} = T_{v_2}^{n+1} - T_{v_1}^{n+1}$$

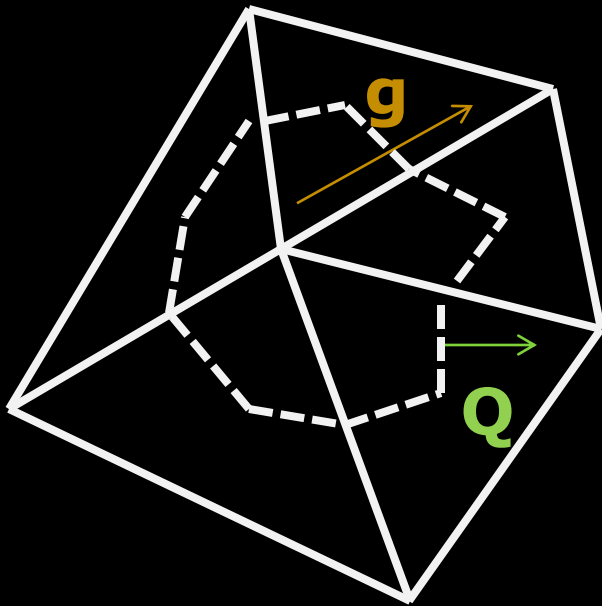
Math: still exact

# Exact Discrete System

$$\begin{aligned} \frac{\partial i}{\partial t} &= -\nabla \cdot \mathbf{q} \\ \mathbf{g} &= \nabla T \end{aligned} \quad \longrightarrow \quad \begin{bmatrix} \mathbf{I} & \mathbf{D} & 0 & 0 \\ 0 & 0 & \mathbf{I} & -\mathbf{G} \end{bmatrix} \begin{pmatrix} I_{\tilde{c}}^{n+1} \\ \bar{Q}_{\tilde{f}} \\ g_e^{n+1} \\ T_v^{n+1} \end{pmatrix} = \begin{pmatrix} I_{\tilde{c}}^n \\ 0 \end{pmatrix}$$

- **Exact**
- **Over-determined**
- **Uncoupled**
- **Unknowns on different meshes**

# Vertex Centered Mimetic



- (C1) Vertex Centered
- (C2) Median dual mesh

$$i \approx \rho C T$$

$$q \approx -k g \quad \Rightarrow \quad Q_{\tilde{f}} \approx -\frac{k\Delta t}{2} (\mathbf{B}^{n+1} g_e^{n+1} + \mathbf{B}^n g_e^n)$$

$$I_{\tilde{c}}^{n+1} \approx M_{\tilde{c}} C_{\tilde{c}} T_v$$

Approximation Part

# Solvable Vertex System

$$\begin{bmatrix} \mathbf{I} & \mathbf{D} & 0 & 0 \\ \mathbf{I} & 0 & 0 & -\mathbf{M}_C \\ 0 & 0 & \mathbf{I} & -\mathbf{G} \\ 0 & \mathbf{I} & \frac{k\Delta t}{2} \mathbf{B}^{n+1} & 0 \end{bmatrix} \begin{pmatrix} I_{\tilde{c}}^{n+1} \\ \bar{Q}_{\tilde{f}} \\ \mathbf{g}_e^{n+1} \\ T_v^{n+1} \end{pmatrix} = \begin{pmatrix} I_{\tilde{v}}^n \\ 0 \\ 0 \\ \frac{k\Delta t}{2} \mathbf{B}^n \mathbf{g}_e^n \end{pmatrix}$$

$$\mathbf{M}_C T_v^{n+1} - I_{\tilde{c}}^n = \mathbf{D} \left( \frac{\Delta t k}{2} \mathbf{B}^{n+1} \right) \mathbf{G} T_v^{n+1} + \mathbf{D} \left( \frac{\Delta t k}{2} \mathbf{B}^n \right) \mathbf{G} T_v^n$$

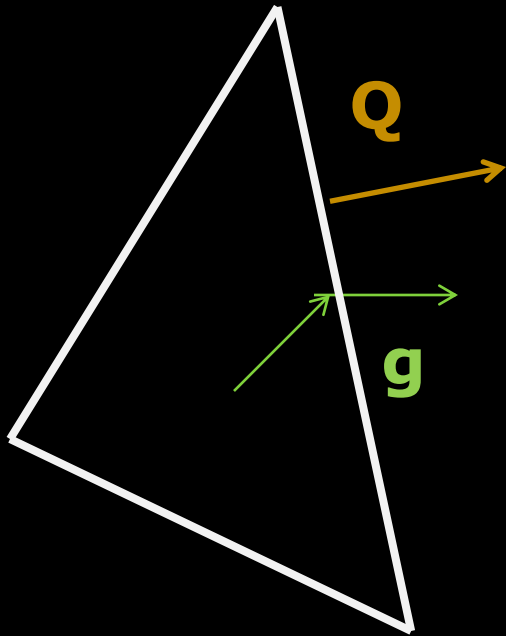
- **Single unknown (T at vertices).**
- **Looks like no dual mesh was used.**
- **B is Hodge \* operator.**



# Other Possibilities

- **Mesh: Use polygons as the primary mesh.**  
(quad/hex meshes, particle methods, SOM)
- **Approx: Use other basis functions for interpolation.**  
(Rational polynomials, Natural Neighbors, Fourier)
- **Exact: Use different exact discretizations**  
(FE, Keller/Priessman Box Schemes)

# Cell Centered Mimetic



- Cell Centered
- Median dual mesh
- Raviart-Thomas for  $Q$

$$i \approx \rho CT$$

$$\mathbf{q} \approx -k\mathbf{g}$$



$$I_c^{n+1} \approx M_c C_c T_{\tilde{c}}$$

$$\frac{\Delta t}{2} (g_{\tilde{e}}^{n+1} + g_{\tilde{e}}^n) \approx -\mathbf{A}_{1/k_c}^{n+1} \overline{Q}_f$$



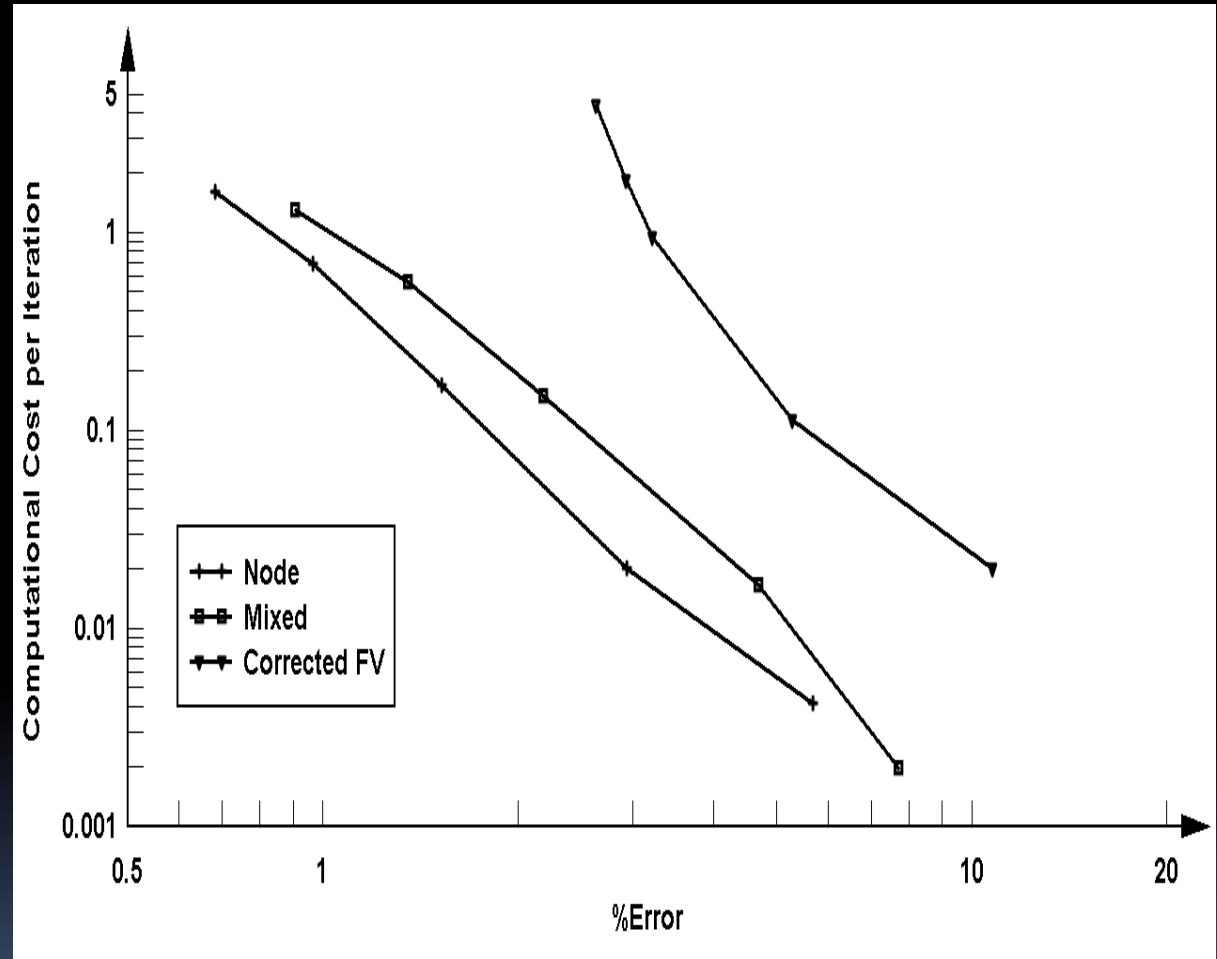
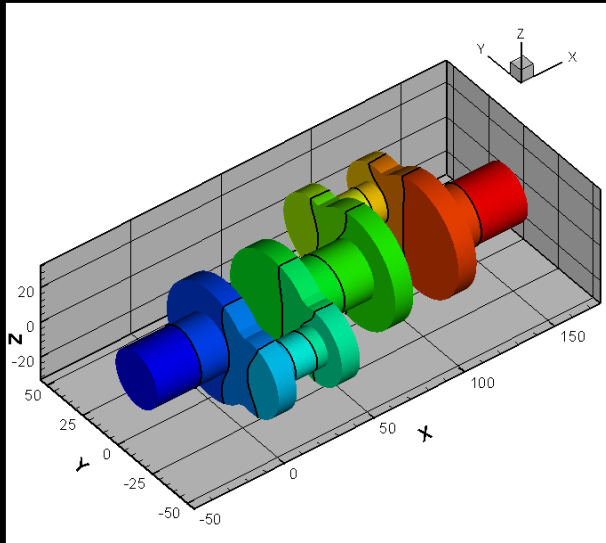
# Cell System

$$\begin{bmatrix} \mathbf{I} & \mathbf{D} & 0 & 0 \\ \mathbf{I} & 0 & 0 & -\mathbf{M}_C \\ 0 & 0 & \mathbf{I} & -\mathbf{G} \\ 0 & \mathbf{A}_{1/k_c}^{n+1} & -\frac{\Delta t}{2}\mathbf{I} & 0 \end{bmatrix} \begin{pmatrix} I_c^{n+1} \\ \bar{Q}_f \\ g_{\tilde{e}}^{n+1} \\ T_{\tilde{c}}^{n+1} \end{pmatrix} = \begin{pmatrix} I_c^n \\ 0 \\ 0 \\ \frac{\Delta t}{2} g_{\tilde{e}}^n \end{pmatrix}$$

$$\begin{bmatrix} \mathbf{M}_C & \mathbf{D} \\ -\mathbf{G} & -\frac{2}{\Delta t} \mathbf{A}_{1/k_c}^{n+1} \end{bmatrix} \begin{pmatrix} T_{\tilde{c}}^{n+1} \\ \bar{Q}_f \end{pmatrix} = \begin{pmatrix} I_c^n \\ g_{\tilde{e}}^n \end{pmatrix}$$

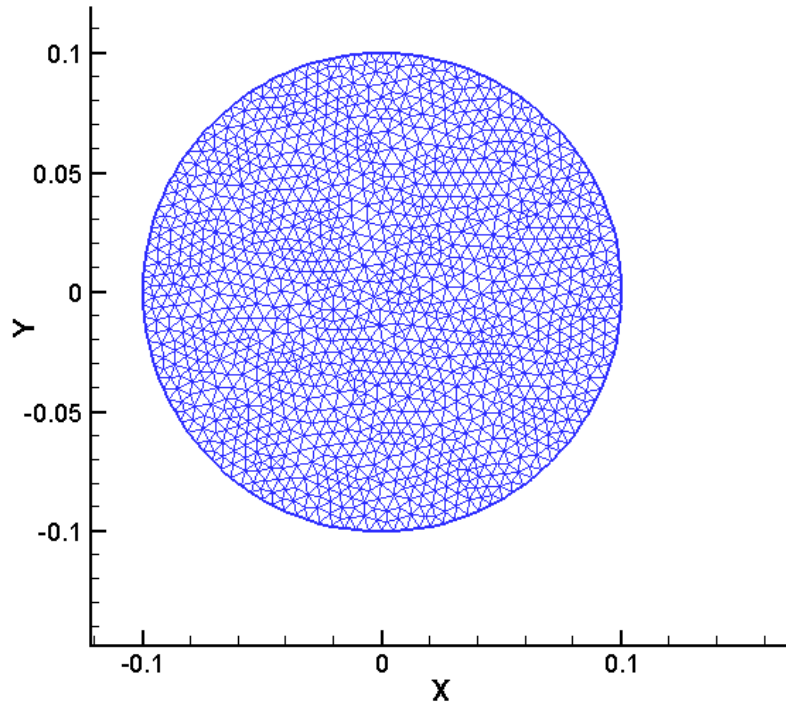
- **Symmetric system of unknowns**
- **Not reducible**

# Comparison

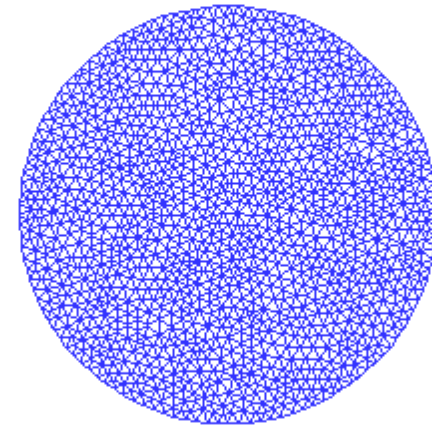


**Mimetic 10x more accurate or 1/10<sup>th</sup> cost.**

# Fluid Example - Droplets

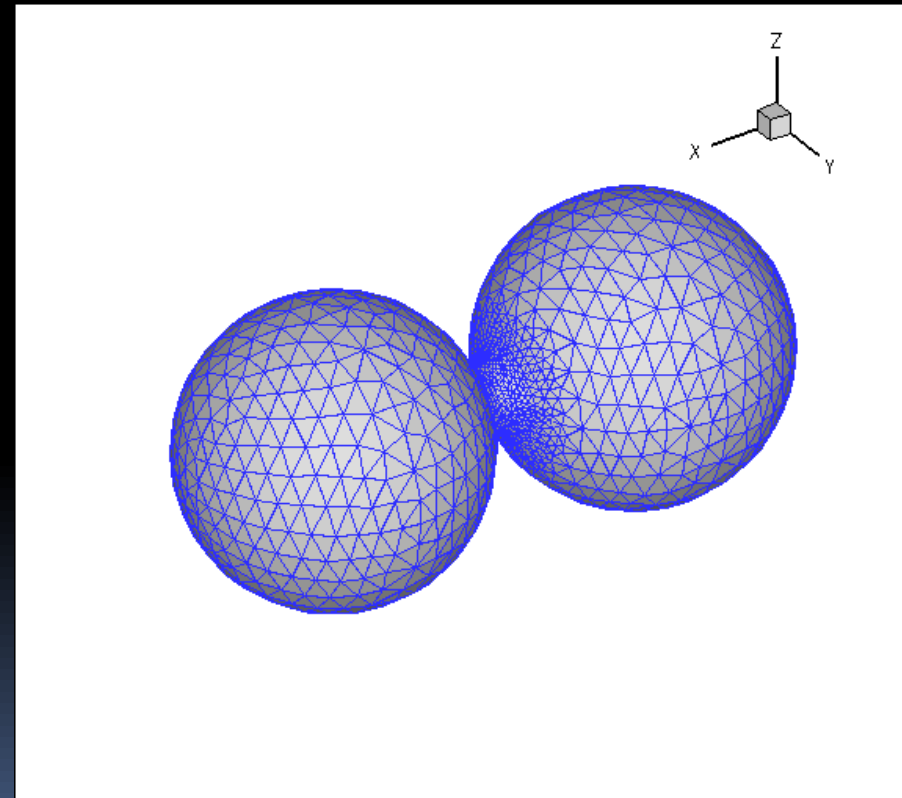
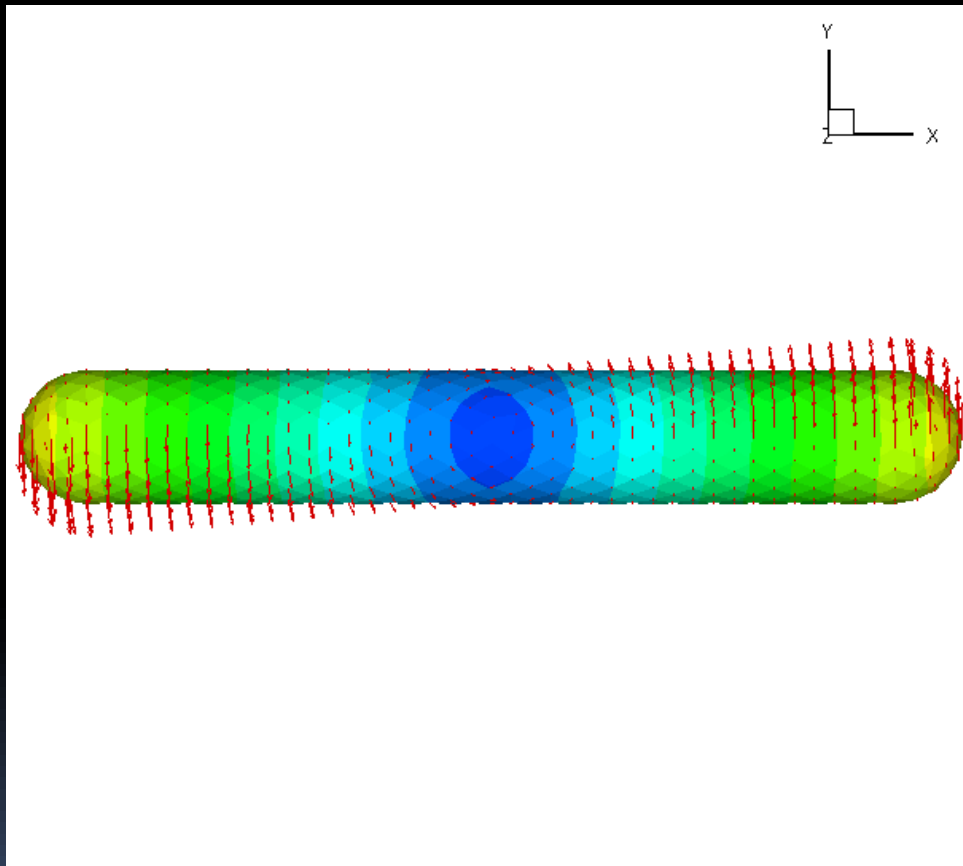


- Moving Mesh
- Unstructured
- Surface Tension

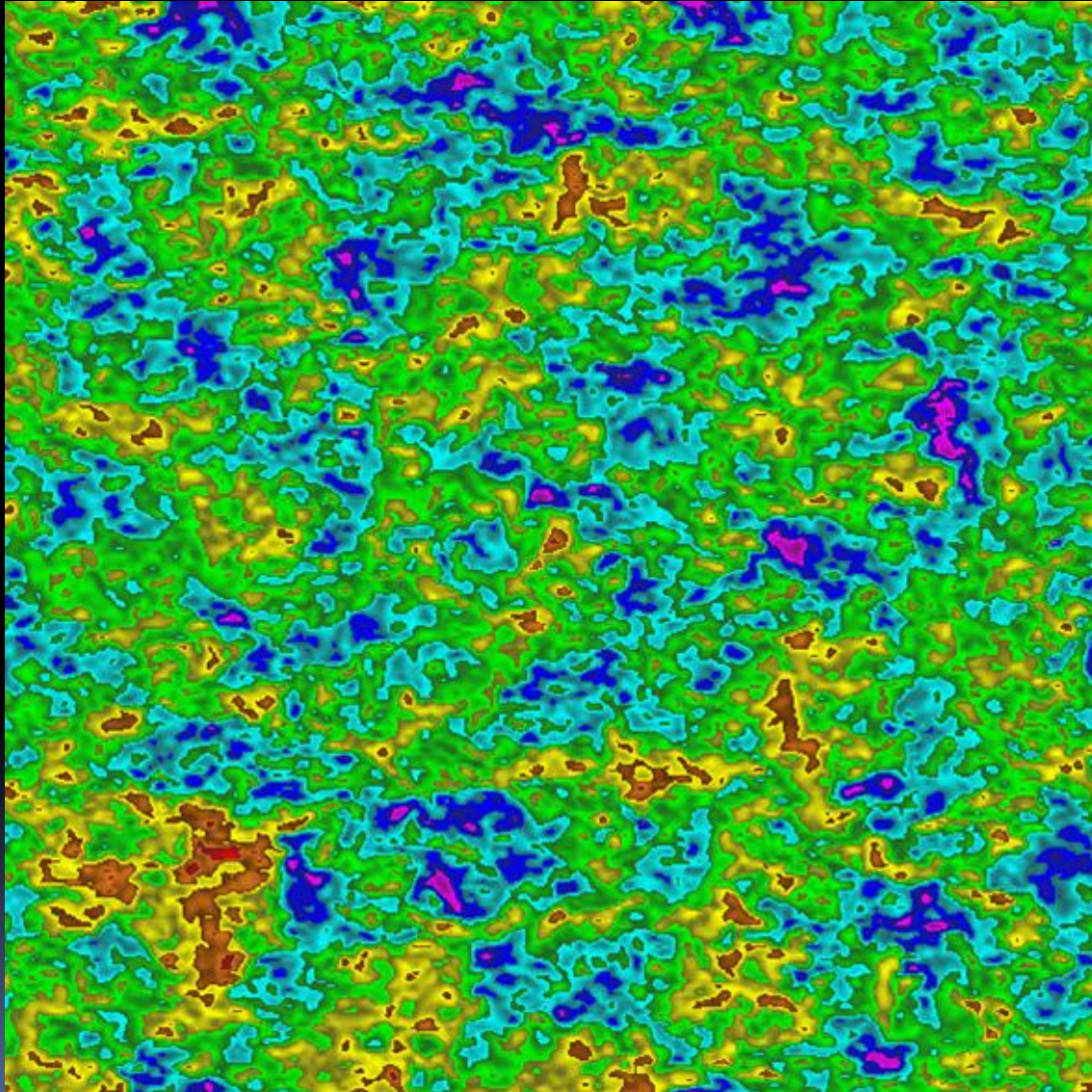


# Example – 3D Droplets

- Moving Mesh
- Large distortions
- Complex BCs



# DNS



# Summary

- **Exact discretization (or 2 step construction) leads to mimetic methods.**
- **Place numerical approximations with the physical approximation.**
- **Discrete Calculus analysis is accessible to all computational scientist. (Gauss' Theorem).**
- **Discrete Calculus Methods are not a type of numerical method. The approach can produce some FV, FE, FD, and other methods.**

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