



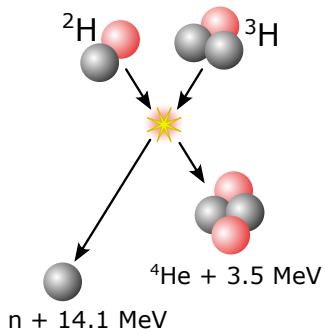
# Discrete adjoint Monte Carlo simulation with reversible random number generators



E. Løvbak

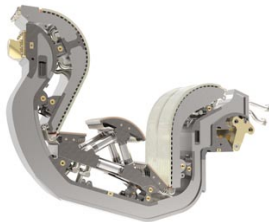
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# Motivation: nuclear fusion in tokamaks



Proton   
Neutron 



$Q \geq 10$  : 50MW  $\rightarrow$  500MW



china eu india japan korea russia usa

## Didactical example: Heat equation in 1D

### ► Macroscopic view

$$\frac{\partial}{\partial t} \mathcal{T}(x, t) - \frac{\partial^2}{\partial x^2} \mathcal{T}(x, t) + u(x) \mathcal{T}(x, t) = 0$$

### ► Particle view

$$\{Q_{p,\tau}\}_{p=1}^P = \left\{ \left[ \begin{array}{c} X_{p,\tau} \\ W_{p,\tau} \end{array} \right] \right\}_{p=1}^P, \quad \tau = 0, \dots, T-1$$

- Diffusion step

$$X_{p,\tau+1} = X_{p,\tau} + \sqrt{2\Delta t} \xi_{p,\tau} \quad \xi_{p,\tau} \sim \mathcal{N}(0, 1)$$

- Reweighting step

$$W_{p,\tau+1} = W_{p,\tau} \exp(-\Delta t \hat{u}(X_{p,\tau+1}))$$

## Optimization problem

### ► Continuous

$$\min_{u(x)} \mathcal{J}(\mathcal{T}(x, t), u(x, t)) = \int_0^\infty \int_0^L \frac{1}{2} \mathcal{T}(x, t)^2 dx dt + \kappa \int_0^L \frac{1}{2} u(x)^2 dx$$

$$\text{subject to } \frac{\partial}{\partial t} \mathcal{T}(x, t) - \frac{\partial^2}{\partial x^2} \mathcal{T}(x, t) + u(x) \mathcal{T}(x, t) = 0,$$

$$\mathcal{T}(x, 0) = \mathcal{T}_0(x), \quad \mathcal{T}(0, t) = \mathcal{T}(L, t)$$

### ► Discrete

$$\mathcal{J}(\mathcal{T}(x, t), u(x, t)) \approx \hat{\mathcal{J}}(\hat{\mathcal{T}}, \hat{u}) = \Delta t \sum_{\tau=0}^T \Delta x \frac{1}{2} \hat{\mathcal{T}}_\tau^\top \hat{\mathcal{T}}_\tau + \nu \Delta x \frac{1}{2} \hat{u}^\top \hat{u}$$

$$\hat{\mathcal{T}}_{\tau, n} = \sum_{p=1}^P \frac{1}{\Delta x} \mathcal{I}_n(X_{p, \tau}) W_{p, \tau}$$

## Adjoint-based optimization

- ▶ PDE-constrained optimization

$$\min_u \mathcal{J}(q, u), \quad \text{subject to} \quad \mathcal{B}(q; u) = 0$$

- ▶ Naive application of chain rule

$$\frac{d\mathcal{J}}{du}(q(u), u) = \frac{\partial \mathcal{J}}{\partial u}(q, u) + \frac{\partial \mathcal{J}}{\partial q}(q, u) \frac{dq}{du}(u)$$

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- ▶ Solution: Lagrangian  $\mathcal{L}(q, q^*, u) = \mathcal{J}(q, u) + (q^*, \mathcal{B}(q, u))$

$$\mathcal{B}(q, u) = 0 \quad \text{State equation}$$

$$\frac{\partial \mathcal{J}^*}{\partial q}(q, u) + \frac{\partial \mathcal{B}^*}{\partial q}(q, u) q^* = 0 \quad \text{Adjoint equation}$$

$$\frac{\partial \mathcal{J}^*}{\partial u}(q, u) + \frac{\partial \mathcal{B}^*}{\partial u}(q, u) q^* = 0 \quad \text{Design equation}$$

## Discrete adjoint

- ▶ Linearized discretization around  $\hat{q}' = \hat{q}(\hat{u})$

$$\hat{\mathcal{B}}(\hat{q}, \hat{u}) \approx \hat{\mathcal{B}}(\hat{q}', \hat{u}) + \frac{\partial \hat{\mathcal{B}}}{\partial \hat{q}}(\hat{q}', \hat{u})(\hat{q} - \hat{q}') = 0$$

is a matrix-vector system



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- ▶ Adjoint equation

$$\frac{\partial \hat{\mathcal{J}}}{\partial \hat{q}}(\hat{q}', \hat{u})^\top + \frac{\partial \hat{\mathcal{B}}}{\partial \hat{q}}(\hat{q}', \hat{u})^\top \hat{q}^* = 0$$



## Adjoint Monte Carlo

- ▶ Simulate particles from final state

$$Q_{p,T}^* = \begin{bmatrix} X_{p,T}^* \\ W_{p,T}^* \end{bmatrix} = -\frac{\partial \hat{\mathcal{J}}^\top}{\partial Q_{p,T}} = -\begin{bmatrix} 0 \\ \Delta t \hat{\mathcal{T}}(X_{p,T}) \end{bmatrix}$$

- ▶ Reverse time stepping

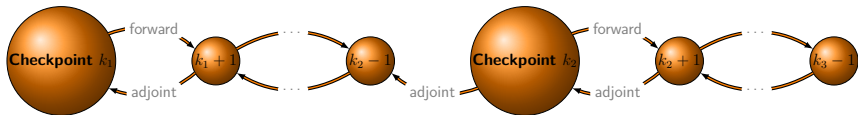
$$\begin{aligned} Q_{p,\tau}^* &= -\frac{\partial B_{p,\tau+1}^\top}{\partial Q_{p,\tau}} Q_{p,\tau+1}^* - \frac{\partial \hat{\mathcal{J}}^\top}{\partial Q_{p,\tau}} \\ &= \begin{bmatrix} X_{p,\tau+1}^* \\ \exp(-\Delta t \hat{u}(X_{p,\tau+1})) W_{p,\tau+1}^* \end{bmatrix} - \begin{bmatrix} 0 \\ \Delta t \hat{\mathcal{T}}(X_{p,\tau}) \end{bmatrix} \end{aligned}$$

- ▶ Note that  $X_{p,\tau}^* = 0$  for all  $\tau$

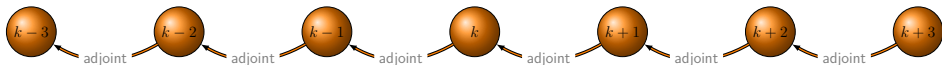
$$W_{p,\tau}^* = \exp(-\Delta t \hat{u}(X_{p,\tau+1})) W_{p,\tau+1}^* - \Delta t \hat{\mathcal{T}}(X_{p,\tau})$$

# Matching forward and adjoint simulations

- ▶ Same paths in forward/backward simulation
- ▶ Challenge:  $P \times T$  large
- ▶ Solutions:
  - Checkpointing: 2 forward simulations + backward simulation



- Generating the paths in reverse (this work)



## Reversing a random number generator

▶ PCG: permuted congruential generator<sup>1</sup>

- internal state  $\zeta_k$  and constant vectors  $a, c, m$

$$\zeta_{k+1} = a\zeta_k + c \pmod{m},$$

- 1-way (permutation) function generates output from  $\zeta_k$
- Passes TestU01 with flying colors

▶ Reversing modular operations  $\rightarrow$  reversed uniform sequence

$$\zeta_k = a^{-1}(\zeta_{k+1} - c) \pmod{m},$$

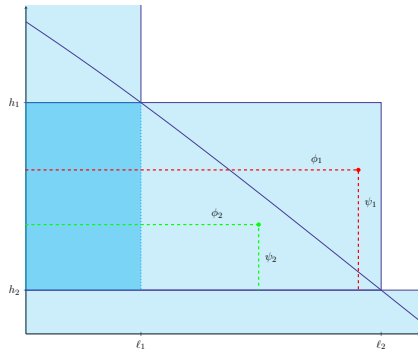
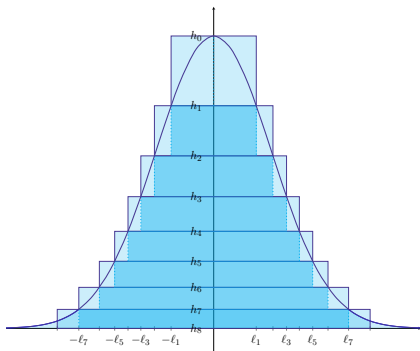
$$a^{-1} \equiv a^{m-2} \pmod{m}$$

▶ Exponential distribution through inverse transform:

$$u \sim \mathcal{U}([0, 1]) \Rightarrow -\lambda \ln(1 - u) \sim \mathcal{E}(\lambda)$$

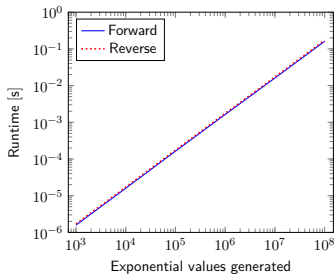
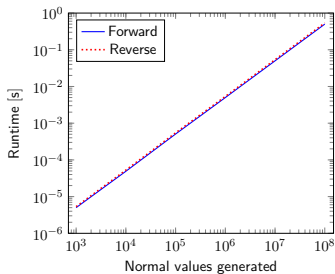
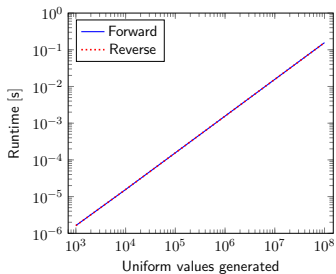
1: M.E. O'Neill, *PCG: A Family of Simple Fast Space-Efficient Statistically Good Algorithms for Random Number Generation*. Technical report HMC-CS-2014-0905, Harvey Mudd College (2014)

# Normal distribution through Ziggurat

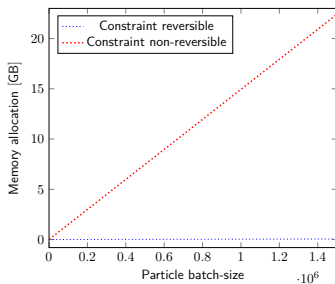
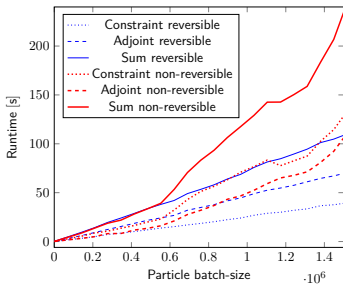
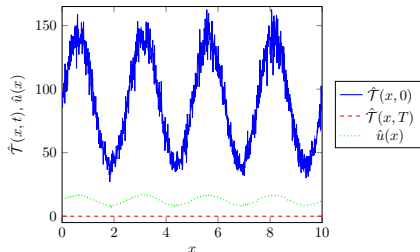


- ▶ Uses either 1, 2 or  $1 + 2n$ ,  $n = 1, 2, \dots$  uniform values
- ▶ How many depends on the first value
- ▶ Solution: Seed second generator

# Generator timings

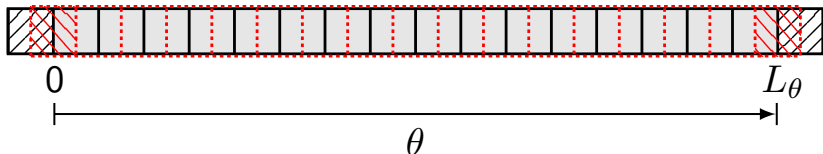


# Optimal cooling





## Fusion example: domain length optimization<sup>2</sup>

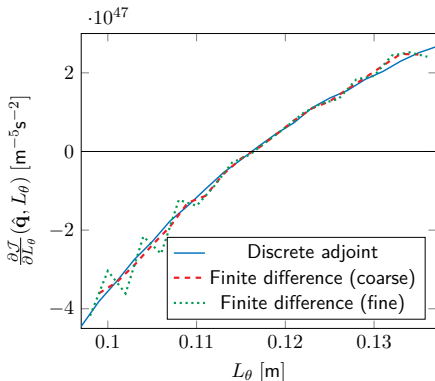
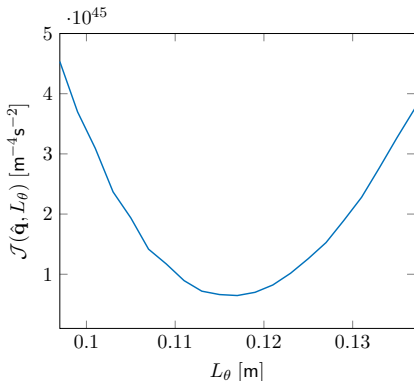


$$\mathcal{J}(q, L_\theta) = \frac{1}{2} \left( n_i b_\theta u_\parallel - \Gamma_d \right)^2 \Big|_{L_\theta} + \frac{\kappa}{2} (L_\theta - L_0)^2$$

$$q = (n_i, u_\parallel, f_n)^\top$$

- $n_i$  plasma density
- $u_\parallel$  plasma velocity
- $f_n$  neutral position-velocity

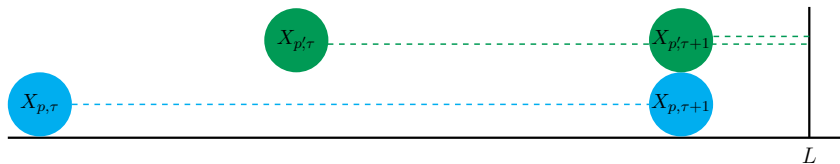
# Preliminary results



	$n_{\mathbf{i}}$ -solver	$u_{  }$ -solver	Neutral solver
State	$3.5 \times 10^{-5}$ s	$4.1 \times 10^{-5}$ s	6.6 s
Adjoint	$3.1 \times 10^{-5}$ s	$3.6 \times 10^{-5}$ s	9.8 s

# Can all simulations be reversed?

- ▶ In theory yes!
- ▶ ... but at what cost?
  - Reflections
  - Collisions in batched simulations
  - ...



Løvnbak, E., Blondeel, F., Lee, A., Vanroye, L., Van Barel, A., Samaey, G., *Reversible random number generation for adjoint Monte Carlo simulation of the heat equation*. Monte Carlo and Quasi-Monte Carlo Methods - MCQMC 2022. Submitted (2023) *arXiv:2302.02778*

Løvnbak, E., *Multilevel and adjoint Monte Carlo methods for plasma edge neutral particle models*. PhD thesis (2023)

## Plasma edge model

- ▶ Plasma  $\Rightarrow$  Finite volume

$$\frac{\partial}{\partial \theta} \left( n_i b_\theta u_{\parallel} \right) = S_{n_i} - K_d n_i$$
$$\frac{\partial}{\partial \theta} \left( m n_i b_\theta u_{\parallel}^2 - \nu_d \frac{\partial u_{\parallel}}{\partial \theta} \right) = S_{u_{\parallel}} - b_\theta \frac{\partial p}{\partial \theta}$$

$$S_\psi = \int f_{\mathbf{n}}(\theta, v) \Psi(\theta, v) dv, \quad \psi \in \{n_i/u_{\parallel}\} \quad \Rightarrow \quad \text{Low-dimensional}$$

- ▶ Neutrals  $\Rightarrow$  Monte Carlo

$$v \frac{\partial}{\partial \theta} f_{\mathbf{n}}(\theta, v) + K_i f_{\mathbf{n}}(\theta, v) = S_{f_{\mathbf{n}}}(n_i, u_{\parallel}) + K_{cx} \int f_{\mathbf{n}}(\theta, v') C(v' \rightarrow v) dv'$$